## Copyright Warning \& Restrictions

The copyright law of the United States (Title 17, United States Code) governs the making of photocopies or other reproductions of copyrighted material.

Under certain conditions specified in the law, libraries and archives are authorized to furnish a photocopy or other reproduction. One of these specified conditions is that the photocopy or reproduction is not to be "used for any purpose other than private study, scholarship, or research." If $a$, user makes a request for, or later uses, a photocopy or reproduction for purposes in excess of "fair use" that user may be liable for copyright infringement,

This institution reserves the right to refuse to accept a copying order if, in its judgment, fulfillment of the order would involve violation of copyright law.

Please Note: The author retains the copyright while the New Jersey Institute of Technology reserves the right to distribute this thesis or dissertation

> Printing note: If you do not wish to print this page, then select "Pages from: first page \# to: last page \#" on the print dialog screen

The Van Houten library has removed some of the personal information and all signatures from the approval page and biographical sketches of theses and dissertations in order to protect the identity of NJIT graduates and faculty.


#### Abstract

\title{ ANALYSIS OF SAFE SEPARATION DISTANCES FROM NATURAL GAS TRANSMISSION PIPELINES } by

\section*{James Stephen Haklar}

Accidental rupture of natural gas transmission pipelines with subsequent ignition of the escaping gas can result in the loss of life and property. In the United States, current means of protecting both the pipeline and the public include the establishment of class locations, in which specific pipeline design, construction and operation requirements must be attained. In addition, distances at which the pipelines can be safely set back from the community, called safe separation distances, have been developed in some European countries (e.g., Great Britain, the Netherlands) through use of risk assessment principles. However, to date there has been no simple, consistent method for determining these distances.

A method for evaluating safe separation distances is proposed herein, in which the point source method for determining heat flux is coupled with relationships for predicting both the mass release rate from the rupture and the flame height of the ignited gas. The method is utilized to develop charts for predicting safe separation distances based on pipeline operating pressure and nominal pipeline diameter. The method is compared to information from both actual pipeline accidents reported upon by the National Transportation Safety Board and from the work of prior researchers utilizing other methodologies. The comparisons reveal that the method proposed in this thesis can produce results that are consistent with the above sets of data tested.




# ANALYSIS OF SAFE SEPARATION DISTANCES <br> FROM <br> NATURAL GAS TRANSMISSION PIPELINES 

by<br>James Stephen Haklar

A Dissertation<br>Submitted to the Faculty of New Jersey Institute of Technology<br>in Partial Fulfillment of the Requirements for the Degree of Doctor of Philosophy<br>Department of Civil and Environmental Engineering

October 1997

Copyright © 1997 by James Stephen Haklar

## ALL RIGHTS RESERVED

## APPROVAL PAGE

# ANALYSIS OF SAFE SEPARATION DISTANCES <br> FROM <br> NATURAL GAS TRANSMISSION PIPELINES 

James Stephen Haklar

Dr. Robert Dresnack, Dissertation Advisor
Professor of Civil and Environmental Engineering, NJIT, Newark, N.J. Date
Dr. Paul C. Chan, Professor of Civil and Environmental Date Engineering, NJIT, Newark, N.J.
Dr. Su Ling Cheng, Professor of CivT Tand Environmental
Engineering, NJIT, Newark, N.J.
Dr. Eygene B. Golu6, Professor of Civil and Envirommental
Engueering, NJIf, Newark, N.I.

| Dr. Robert P. Kirchner, Professor of Mechanical |
| :--- |
| Engineering, NIIT, Newark, N.J. |

## BIOGRAPHICAL SKETCH

Author: James Stephen Haklar
Degree: Doctor of Philosophy
Date: October 1997

## Date of Birth:

## Place of Birth:

## Undergraduate and Graduate Education:

- Doctor of Philosophy in Civil Engineering, New Jersey Institute of Technology, Newark, NJ, 1997
- Master of Science in Environmental Engineering, New Jersey Institute of Technology, Newark, NJ, 1989
- Bachelor of Science in Chemical Engineering, New Jersey Institute of Technology, Newark, NJ, 1984

Major: Civil Engineering

## Professional Experience:

- United States Environmental Protection Agency, 290 Broadway, New York, NY, 1985 - Present

Title: Environmental Engineer

To my wife, my parents, and my sister

## ACKNOWLEDGMENT

The author wishes to express his sincerest appreciation to his advisor, Professor Robert Dresnack, for his support, guidance and friendship. Special thanks to Professors Eugene B. Golub, Paul C. Chan, Su Ling Cheng and Robert P. Kirchner for participating as committee members.

The author is grateful to the staff of the National Transportation Safety Board, the Transportation Safety Board of Canada, the Canadian National Energy Board, the United States Environmental Protection Agency and the municipal offices of the Township of Edison, New Jersey for all of the information and guidance concerning natural gas pipeline accidents.

Lastly, special thanks to Raymond Basso, Janet Feldstein, Lisa Jackson and Carole Petersen for their patience and encouragement during the development of this thesis.

## TABLE OF CONTENTS

Chapter Page
I INTRODUCTION ..... 1
2 LITERATURE REVIEW ..... 6
2.1 Natural Gas and its Properties ..... 7
2.2 Compressible Fluid Flow ..... 9
2.2.1 Choked Conditions ..... 9
2.2.2 Isothermal Pipe Flow ..... 11
2.2.3 Adiabatic Pipe Flow ..... 12
2.3 Prior Research Pertaining to Pipeline Blowdown ..... 12
2.3.1 Wilson ..... 12
2.3.2 Groves, Bishnoi and Wallbridge ..... 14
2.3.3 Fannelop and Ryhming ..... 15
2.3.4 Flatt ..... 17
2.3.5 Lang and Fannelop ..... 18
2.3.6 Ryhming ..... 18
2.3.7 Picard and Bishnoi ..... 19
2.3.8 Olorunmaiye and Imide ..... 20
2.3.9 Other Work ..... 20

## TABLE OF CONTENTS (Continued)

Chapter Page
2.4 Prior Research Pertaining to Turbulent Jet Flames and Flame Geometry ..... 20
2.4.1 Turbulent Jet Flames ..... 20
2.4.2 Flame Geometry ..... 21
2.4.3 Heat Transfer Models ..... 24
2.4.3.1 Point Source Model ..... 24
2.4.3.2 Line Source Model ..... 25
2.4.3.3 Tilted Cylinder Model ..... 26
2.5 Prior Research in the Field of Hazard Analysis ..... 26
2.5.1 Typical Models ..... 26
2.5.2 Research Performed by the European
Gas Pipeline Incident Data Group ..... 28
2.5.3 Research Performed by the British
Health and Safery Executive ..... 29
2.5.4 Research Performed in Canada ..... 30
2.5.5 Research Performed by NJIT ..... 32
2.6 Conclusions from Literature Review ..... 32
3 OBIECTIVE OF THE DISSERTATION AND INTRODUCTION TO THE DISSERTATION RATIONALE ..... 36
4 DEVELOPMENT OF METHOD TO DETERMINE SAFE SEPARATION DISTANCES ..... 41

# TABLE OF CONTENTS <br> (Continued) 

Chapter Page
4.1 Calculation of Burn Radii ..... 41
4.1.1 Use of the Point Source Method ..... 42
4.1.2 Heat Flux Values ..... 44
4.1.3 Determination of Gas Flow Rate ..... 46
4.1.4 Determination of Flame Height ..... 50
4.1.5 Construction of Charts to Predict Safe Separation Distances ..... 53
5 COMPARISON OF METHOD TO PIPELINE ACCIDENT DATA AND TO PREVIOUS RESEARCH ..... 55
5.1 Comparison to Pipeline Accident Data ..... 56
5.1.1 PAR 95-01 ..... 56
5.1.2 PAR 87-01 ..... 61
5.1.3 Other Data ..... 66
5.2 Comparison to Separation Distances Developed through Hazard Analysis ..... 70
5.2.1 Separation Distances Imposed in the Dutch Regulations ..... 70
5.2.2 Separation Distances Determined by the British Health and Safety Executive ..... 73

## TABLE OF CONTENTS (Continued)

Chapter Page
6 RESULTS AND DISCUSSIONS ..... 76
6.1 Effect of Variation in Heat Flux on Burn Radius ..... 77
6.2 Effect of Variation in Atmospheric Transmissivity on Burn Radius ..... 81
6.3 Effect of Variation in Wind Speed and Direction on Burn Radius ..... 85
7 CONCLUSIONS AND RECOMMENDATIONS ..... 89
APPENDIX A - HEAT FLUX LEVELS ..... 91
APPENDIX B - BURN RADIUS CHARTS ..... 100
APPENDIX C - FIGURES AND TABLES FOR CHAPTER 5 ..... 129
APPENDIX D - VARIATION OF BURN RADIUS INFORMATION ..... 137
APPENDIX E - CANADIAN HEAT FLUX INFORMATION ..... 147
REFERENCES ..... 161

## LIST OF TABLES

Table Page
2.1 Composition of Natural Gas ..... 7
2.2 Composition of Sour Gas ..... 9
4.1 Heat Flux Values ..... 45
4.2 Pipeline Rupture Parameters ..... 47
4.3 Flow Rate Parameters ..... 48
4.4 Flame Height Data ..... 51
5.1 Heat Flux Experienced at Lancaster, Kentucky Accident ..... 65
5.2 Comparison of Health and Safety Executive Results with Equation 4-12 ..... 73
5.3 Comparison of all Natural Gas Pipelines ..... 74
6.1 Pipeline Conditions ..... 77
6.2 Burn Radii at Various Heat Fluxes ..... 78
6.3 Variation of Heat Flux Values ..... 79
6.4 Variation of Heat Flux Values (Contribution of Solar Heat Flux) ..... 80
6.5 Atmospheric Transmissivities ..... 82
6.6 Wind Speeds for Several U.S. Cities ..... 86
A. 1 Summary of Heat Flux Levels ..... 92
C. 1 Comparison of Building Distances to Safe Separation Distances ..... 134
C. 2 Ratios of Building and Separation Distances at Various Diameters (Constant Pressure) ..... 135

## LIST OF TABLES

## (Continued)

Table Page
C. 3 Ratios of Building and Separation Distances at Various Pressures (Constant Diameter) ..... 136
D. I Variation of Burn Radius with Atmospheric Transmissivity ..... 139
D. 2 Percent Difference Between Burn Radii Obtained Using Assumed Value of $\tau$ and Burn Radii Determined Using Maximum and Minimum Values of $\tau$ ..... 140
D. 3 Variation of Burn Radii for Combined Variations in Heat Flux and Atmospheric Transmissivity ..... 141
D. 4 Percent Difference in Burn Radii for Combined Variations in Heat Flux and Atmospheric Transmissivity ..... 142
D. 5 Variation of Burn Radius with Wind Direction and Speed $24^{\prime \prime}$ Diameter Pipeline, $\mathrm{P}=1,000 \mathrm{psig}$ ..... 143
D. 6 Percent Difference Between Equation 4-12 and the Canadian Values $24^{\prime \prime}$ Diameter Pipeline, $P=1,000$ psig ..... 144
D. 7 Variation of Burn Radius with Wind Direction and Speed
$36^{\prime \prime}$ Diameter Pipeline, $\mathrm{P}=1,000 \mathrm{psig}$ ..... 145
D. 8 Percent Difference Between Equation 4-12 and the Canadian Values 36" Diameter Pipeline, $\mathrm{P}=1,000 \mathrm{psig}$ ..... 146

## LIST OF FIGURES

Figure Page
2.1 Flame Shape with a Crosswind ..... facing 22
2.2 Jet Flame Shape Parameters ..... facing 23
2.3 Conceptual Illustration of the Point Source Model ..... facing 24
2.4 Line Source Model ..... facing 25
4.1 Point Source Method Applied to Pipeline Rupture ..... facing 42
4.2 Example of a Burn Radius Chart for a $36^{\prime \prime}$
Diameter Pipeline, $K=9,985$ BTU/hr sq. ft. ..... facing 54
5.1 Location of Edison, New Jersey Accident ..... facing 57
5.2 Location of Lancaster, Kentucky Accident ..... facing 61
B. 1 Burn Radius Chart for $14^{\prime \prime}$ Diameter Pipeline, $\mathrm{K}=3,962 \mathrm{Btu} / \mathrm{hr} \mathrm{sq}$. ft. ..... 101
B. 2 Burn Radius Chart for $16^{\prime \prime}$ Diameter Pipeline, $\mathrm{K}=3,962 \mathrm{Btu} / \mathrm{hr} \mathrm{sq}. \mathrm{ft}$. ..... 102
B. 3 Burn Radius Chart for 18" Diameter Pipeline, $K=3,962 \mathrm{Btu} / \mathrm{hr} \mathrm{sq}$. ft. ..... 103
B. 4 Burn Radius Chart for 20 " Diameter Pipeline, $\mathrm{K}=3,962 \mathrm{Btu} / \mathrm{hr} \mathrm{sq}$. ft. ..... 104
B. 5 Burn Radius Chart for 24" Diameter Pipeline, $\mathrm{K}=3,962 \mathrm{Btu} / \mathrm{hr} \mathrm{sq}$. fi. ..... 105
B. 6 Burn Radius Chart for 30" Diameter Pipeline, $\mathrm{K}=3,962$ Btu/hr sq. ft. ..... 106
B. 7 Burn Radius Chart for 36" Diameter Pipeline, $\mathrm{K}=3,962$ Btwhr sq. ft. ..... 107

## LIST OF FIGURES <br> (Continued)

Figure Page
B. 8 Burn Radius Chart for $14^{\prime \prime}$ Diameter Pipeline, $\mathrm{K}=6,340 \mathrm{Btu} / \mathrm{hr} \mathrm{sq} . \mathrm{ft}$. ..... 108
B. 9 Burn Radius Chart for $16^{\prime \prime}$ Diameter Pipeline, $\mathrm{K}=6,340 \mathrm{Btu} / \mathrm{hr} \mathrm{sq} . \mathrm{ft}$. ..... 109
B. 10 Burn Radius Chart for $18^{\prime \prime}$ Diameter Pipeline, $K=6,340 \mathrm{Btu} / \mathrm{hr} \mathrm{sq} . \mathrm{ft}$. ..... 110
B. 11 Burn Radius Chart for $20^{\prime \prime}$ Diameter Pipeline, $\mathrm{K}=6,340 \mathrm{Btu} / \mathrm{hr} \mathrm{sq} . \mathrm{ft}$. ..... 111
B. 12 Burn Radius Chart for $24^{\prime \prime}$ Diameter Pipeline, $\mathrm{K}=6,340 \mathrm{Btu} / \mathrm{hr} \mathrm{sq} . \mathrm{ft}$. ..... 112
B. 13 Burn Radius Chart for 30" Diameter Pipeline, $\mathrm{K}=6,340 \mathrm{Btu} / \mathrm{hr} \mathrm{sq} . \mathrm{ft}$. ..... 113
B. 14 Burn Radius Chart for $36^{\prime \prime}$ Diameter Pipeline, $K=6,340 \mathrm{Btu} / \mathrm{hr} \mathrm{sq} . \mathrm{ft}$. ..... 114
B. 15 Burn Radius Chart for $14^{\prime \prime}$ Diameter Pipeline, $\mathrm{K}=9,510 \mathrm{Btu} / \mathrm{hr} \mathrm{sq} . \mathrm{ft}$. ..... 115
B. 16 Burn Radius Chart for $16^{\prime \prime}$ Diameter Pipeline, $K=9,510 \mathrm{Btu} / \mathrm{hr} \mathrm{sq} . \mathrm{ft}$. ..... 116
B. 17 Burn Radius Chart for $18^{\prime \prime}$ Diameter Pipeline, $K=9,510 \mathrm{Btu} / \mathrm{hr} \mathrm{sq} . \mathrm{ft}$. ..... 117
B. 18 Burn Radius Chart for 20" Diameter Pipeline, $\mathrm{K}=9,510 \mathrm{Btu} / \mathrm{hr} \mathrm{sq} . \mathrm{ft}$. ..... 118
B. 19 Burn Radius Chart for $24^{\prime \prime}$ Diameter Pipeline,
$K=9,510 \mathrm{Btu} / \mathrm{hr} \mathrm{sq} . \mathrm{ft}$. ..... 119
B. 20 Burn Radius Chart for $30^{\prime \prime}$ Diameter Pipeline, $\mathrm{K}=9,510 \mathrm{Btu} / \mathrm{hr} \mathrm{sq} . \mathrm{ft}$. ..... 120

## LIST OF FIGURES (Continued)

FigureB. 21 Burn Radius Chart for 36" Diameter Pipeline,$K=9,510 \mathrm{Btu} / \mathrm{hr} \mathrm{sq} . \mathrm{ft}$.121
B. 22 Burn Radius Chart for $14^{\prime \prime}$ Diameter Pipeline, $K=9,985 \mathrm{Btu} / \mathrm{hr} \mathrm{sq} . \mathrm{ft}$. ..... 122
B. 23 Burn Radius Chart for $16^{\prime \prime}$ Diameter Pipeline, $K=9,985 \mathrm{Btu} / \mathrm{hr} \mathrm{sq} . \mathrm{ft}$. ..... 123
B. 24 Burn Radius Chart for $18^{\prime \prime}$ Diameter Pipeline, $\mathrm{K}=9,985 \mathrm{Btu} / \mathrm{hr} \mathrm{sq} . \mathrm{ft}$. ..... 124
B. 25 Burn Radius Chart for $20^{\prime \prime}$ Diameter Pipeline, $K=9,985 \mathrm{Btu} / \mathrm{hr}$ sq. ft. ..... 125
B. 26 Burn Radius Chart for $24^{\prime \prime}$ Diameter Pipeline, $K=9,985 \mathrm{Btu} / \mathrm{hr} \mathrm{sq} . \mathrm{ft}$. ..... 126
B. 27 Burn Radius Chart for $30^{\prime \prime}$ Diameter Pipeline, $K=9,985 \mathrm{Btu} / \mathrm{hr} \mathrm{sq} . \mathrm{ft}$. ..... 127
B. 28 Burn Radius Chart for $36^{\prime \prime}$ Diameter Pipeline, $K=9,985 \mathrm{Btu} / \mathrm{hr} \mathrm{sq} . \mathrm{ft}$. ..... 128
C. 1 Detailed View of Edison Accident ..... 130
C. 2 Identification of Apartment Buildings ..... 131
C. 3 Illustration of Apartments Used for Development of an Appropriate Scale ..... 132
C. 4 Comparison of Flux Levels Experienced Versus
Levels for Blistering ..... 133

## LIST OF ABBREVIATIONS AND SYMBOLS

| Abbreviations |  |
| :---: | :---: |
| A | cross-sectional area of pipeline |
| $\mathrm{A}_{2}$ | area of minimum cross-section of a puncture or nozzle |
| a | ambient air condition |
| ALOHA | Aeral Locations of Hazardous Atmospheres |
| API | American Petroleum Institute |
| atm | atmosphere |
| BLEVE | Boiling Liquid Expanding Vapor Explosion |
| BR | burn radius |
| Btu | British thermal unit |
| - |  |
| $C_{L}$ | lean limit concentration |
| c | critical conditions |
| CCPS | Center for Chemical Process Safety |
| CFR | Code of Federal Regulations |
| cm | centimeter |
| $\mathrm{C}_{\mathrm{p}}$ | heat capacity at constant pressure |
| $\mathrm{C}_{\mathrm{T}}$ | mol fraction of nozzle fluid in unreacted stoichiometric mixture |
| C.W. | crosswind |
| $\mathrm{C}_{V}$ | heat capacity at constant volume |
| $\mathrm{D}_{\text {S }}$ | burner source diameter |
| D | diameter, distance from flame center to observer |
| d | pipeline diameter |
| $\mathrm{d}_{\mathrm{j}}$ | jet diameter, diameter of flare stack |
| D.W. | downwind |
| E | longitudinal joint factor, surface emissive power of a flame |
| $E_{b b}$ | equivalent blackbody emissive power |
| EGIG | European Gas Pipeline Incident Data Group |
| F | design factor, geometric view factor, fraction of total heat radiated |
| $f$ | friction factor |
| ${ }^{\circ} \mathrm{F}$ | degrees Fahrenheit |
| FLACS | Flame Acceleration Simulator |
| ft | foot |
| H | flame height |
| h | hour, height of flare stack |
| hr | hour |
| HACS | Hazard Assessment Computer System |
| HSE | Health and Safety Executive |
| I | thermal radiation intensity |

# LIST OF ABBREVIATIONS AND SYMBOLS (Continued) 

| Abbreviations |  |
| :---: | :---: |
| j | discharge condition |
| K | degrees Kelvin, radiant heat flux, constant |
| k | gas specific heat ratio, extinction coefficient |
| kPa | kilopascal |
| kW | kilowatt |
| L | length |
| $L_{\text {B }}$ | distance from burner exit to flame tip |
| $L_{\text {L }}$ V | vertical length of flame tip from the plane of the burner |
| 1 | distance between upstream and downstream locations |
| 1b | pound |
| M | Mach number, mass |
| $M_{1}$ | upstream Mach number |
| $\mathrm{M}_{2}$ | downstream Mach number |
| $\mathrm{Ma}_{\text {a }}$ | molecular weight of air |
| $\mathrm{M}_{\mathrm{f}}$ | molecular weight of fuel |
| $M_{N}$ | molecular weight of nozzle fluid |
| $\mathrm{M}_{\mathrm{S}}$ | molecular weight of surrounding fluid |
| m* | mass flow rate |
| $\mathrm{m}^{*} \mathrm{e}$ | mass release rate of gas at time $=1$ |
| $\mathrm{m}^{*}{ }_{\text {oe }}$ | initial mass release rate |
| m | mete: |
| Mcf | thousand cubic feet |
| MIACC | Major Industrial Co-ordinating Council of Canada |
| mm | millimeter |
| NFPA | National Fire Protection Association |
| NGL | natural gas liquids |
| NJIT | New Jersey Institute of Technology |
| NPS | nominal pipe size |
| NTSB | National Transportation Safety Board |
| P | design pressure, incident operating pressure |
| P' | gauge operating pressure |
| $P_{1}$ | absolute pressure near the opening |
| $p_{1}$ | absolute vessel pressure, upstream pressure, operating pressure at rupture |
| $\mathrm{P}_{2}$ | absolute pressure in the minimum cross-section, downstream pressure |
| PAR | pipeline accident report |
| psi | pounds per square inch |
| psia | pounds per square inch absolute |

## LIST OF ABBREVIATIONS AND SYMBOLS (Continued)

| Abbreviations |  |
| :---: | :---: |
| psig | pounds per square inch gauge |
| Q | total heat content of flared gas, volumetric gas flow rate |
| q | radiant heat flux, thermal radiation at receptor |
| q" | incident radiant flux |
| $\mathrm{q}_{2}$ | average heat flux |
| - | relative humidity |
| R | gas constant, velocity ratio equal to crosswind speed (U)/ jet velocity ( $\mathrm{U}_{\mathrm{j}}$ ) |
| ${ }^{0} \mathrm{R}$ | degrees Rankine |
| ROW | right-of-way |
| RPA | Response Planning Area |
| S | yield strength, curvilinear coordinate |
| - |  |
| $S_{0}$ | curvilinear distance |
| s, sec. | second |
| scf | standard cubic foot |
| SFPE | Society of Fire Protection Engineers |
| SAFEMODE | Safety Assessment for Effective Management of Dangerous Events |
| T | temperature derating factor, temperature |
| $\mathrm{T}_{1}$ | gas temperature |
| $\mathrm{T}_{0}$ | seawater temperature |
| $T_{F}$ | adiabatic flame temperature |
| $\mathrm{T}_{\mathrm{N}}$ | temperature of nozzle fluid |
| t | wall thickness, time, end of flame axis |
| ${ }^{1}$ | duration of exposure |
| TSB | Transportation Safety Board (of Canada) |
| 4 | gas velocity, velocity |
| $\mathrm{u}_{j}$ | velocity of gas jet |
| $\mathrm{u}_{1}$ | velocity of wind |
| U.W. | upwind |
| V' | volumetric gas flow rate |
| $\mathrm{V}_{\text {AVG }}$ | average gas flow rate |
| $V_{\text {CALC }}$ | maximum gas flow rate |
| $W_{0}$ | total mass of gas in pipeline |
| W | flame base width |
| $W_{2}$ | flame tip width |
| WHAZAN | World Bank Hazard Analysis |
| X | $x$-coordinate of flame tip |

# LIST OF ABBREVIATIONS AND SYMBOLS (Continued) 

| Abbreviations |  |
| :---: | :---: |
| - |  |
| x | $x$-coordinate of receptor |
| $\bar{X}_{0}$ | dimensionless coordinate |
| yr | year |
| Z | $z$-coordinate of flame tip |
| z | cross-stream distance |
| - |  |
| z | z-coordinate of receptor |
| $\bar{Z}_{0}$ | dimensionless rise |

## Symbols

| $\alpha$ | $\mathrm{W}_{0} /\left[/ \mathrm{m}^{*}{ }_{\text {oe }}\right]$, angle subtended by flame with respect to vertical |
| :---: | :---: |
| $\alpha_{\text {B }}$ | angle subtended by burner tip and tip of flame with respect to vertical |
| $\alpha_{T}$ | mols reactants/mols products for the stoichiometric mixture |
| $\beta$ | final release time constant |
| $d$ | differential (i.e., derivative) operator |
| d | partial differential (i.e., partial derivative) operator |
| $\theta$ | angle between normal to surface receiving thermal radiation and line of sight from flame center |
| $\rho$ | gas density |
| $\rho_{1}$ | density of gas inside vessel |
| $\rho_{a}$ | density of air |
| $\rho_{j}$ | density of jet |
|  | atmospheric transmissivity |

## LIST OF CONVERSIONS

## Area

1 Acre $=43,560 \mathrm{ft}^{2}$

## Heat Transfer

```
\(1 \mathrm{Btu} / \mathrm{hft}^{2}=3.1546 \mathrm{~W} / \mathrm{m}^{2}\)
\(1 \mathrm{Btu} / \mathrm{h}=0.29307 \mathrm{~W}\)
1 calorie \(/ \mathrm{h}=1.1622 \times 10^{-3} \mathrm{~W}\)
```


## Length

$$
\begin{aligned}
& 1 \mathrm{in}=2.540 \mathrm{~cm} \\
& 100 \mathrm{~cm}=1 \mathrm{~m} \\
& 1 \mathrm{~m}=3.2808 \mathrm{ft}=39.37 \mathrm{in} . \\
& 5,280 \mathrm{ft}=1 \mathrm{mile}
\end{aligned}
$$

## Pressure

$1 \mathrm{~atm}=14.696 \mathrm{psia}=1.01325 \mathrm{bar}$
1 bar $=1 \times 10^{5} \mathrm{~Pa}$
gauge pressure $=$ absolute pressure -1 atm

## Temperature

$T$ (Celsius) $=\mathrm{K}-273.15$
$1.8^{\circ} \mathrm{F}=\mathrm{K}=1^{\circ} \mathrm{C}$

Note: All conversions with the exception of the conversion for Area are from Geankoplis, C.J., Transport Processes and Unit Operations. Boston, MA: Allyn and Bacon, Inc. (1983). The conversion for Area is from Wehster's Ninth New Collegiate Dictionary. Springfield, MA: Merriam Webster, Inc. (1991).

## CHAPTER I

## INTRODUCTION

Natural gas is most commonly transported between collection and distribution points through a series of interstate pipelines called transmission pipelines. According to the Transportation Research Board of the National Research Council, there are approximately 271,000 miles of natural gas transmission lines in the United States. The lines are typically constructed of steel, operate between 400 and 1,500 pounds per square inch ( psi ) of pressure and are buried under 2.5 feet to 3 feet or more of cover (1).

Increased development of formerly sparsely populated areas has resulted in instances of encroachment of transmission pipeline rights-of-way (ROWs). Accidental ruptures of these pipelines with subsequent ignition of the escaping gas can result in the loss of life and property near these lines.

A description of the effects of such accidents can be found in the pipeline accident reports prepared by the National Transportation Safety Board (NTSB). The NTSB is an independent Federal agency that investigates pipeline accidents occurring in the United States. It is worth noting that the Canadian government has established a similar agency, called the Transportation Safety Board of Canada (TSB). The TSB's functions are similar to those of the NTSB, in that the NTSB also investigates transportation accidents.

One example of the destructive effects of pipeline ruptures was the rupture and resultant fire occurring at Beaumont Kentucky, on April 27, 1985 involving a 30-inch diameter natural gas line operating at 992 pounds per square inch gauge (psig). This accident killed 5 persons, injured 3 and, excluding the pipeline, resulted in one million
dollars worth of damage (2). More recently, a 36 -inch pipeline operating at a pressure of 970 psig exploded in Edison, New Jersey'. This accident occurred on March 23, 1994 and severely affected a nearby apartment complex. The apartment complex sustained $\$ 12.4$ million in damages, which includes the loss of eight apartment buildings, severe damage to six buildings and minor damage to several other buildings (3). These two examples serve to illustrate the importance of determining the proximity at which pipelines can be safely sited near a community. These distances are called "safe separation distances."

In the United States, pipelines are regulated by the U.S. Department of Transportation pursuant to the statutory authority of the Natural Gas Pipeline Safety Act of 1968. The implementing regulations can be found in Parts 191 and 192 of the Code of Federal Regulations (CFR). The regulations establish a series of four class locations, within which specific design, construction and operational requirements must be met. Class locations extend 220 yards on either side of the centerline of any 1 mile length of pipeline, and are differentiated as follows (4):

Class 1 Location: Has 10 or less buildings intended for human occupancy.
Class 2 Location: Has more than 10 but less than 46 buildings intended for human occupancy.

Class 3 Location: Has 46 or more buildings intended for human occupancy; or an area where the pipeline lies within 100 yards of either a building or a small, well defined outside area (such as a playground, recreation area, outdoor theater, or other place of public assembly) that is occupied by 20 or more persons on at least 5 days a week for 10 weeks in any 12 month period.

Class 4 Location: Where buildings with four or more stories above ground are prevalent.

It should be noted that each separate living unit in a multiple unit building is counted as a separate building intended for human occupancy.

As indicated above, each class location has specific requirements that pertain to the design, construction and operation of a pipeline. For example, the regulations specify that the design pressure for steel pipe be determined in accordance with the following equation:

## $\mathrm{P}=(2 \mathrm{St} / \mathrm{D}) \times \mathrm{FxExT}$

In the above equation, $P$ represents the design pressure ( psig ); S is the yield strength ( psi ); D is the outside diameter (inches); t is the wall thickness (inches); and $\mathrm{F}, \mathrm{E}$ and T are, respectively, the design factor, longitudinal joint factor and temperature derating factor (all dimensionless). As indicated in the regulations (5), the pressure design factors for steel pipe can vary from 0.72 to 0.40 as the Class Location varies from 1 to 4 . The United States regulations, however, do not specify required separation distances between a natural gas pipeline and buildings in proximity to the pipeline.

Other countries have likewise established regulations for natural gas pipelines (5). Canada has developed a four class location system which is almost identical to the system used in the United States. The German regulations establish separation distances which are dependent on the diameter of the pipe. However, the regulations indicate that the right of way "is to protect the pipeline" (as opposed to protecting the community). Australia's regulations provide for six class locations, but (as with the United States and Canadian regulations) do not specify minimum separation distances between pipelines
and neighboring structures. While the Japanese regulations specify minimum separation distance requirements from select receptors (such as schools, theaters, train stations, etc.), these requirements apply to above-ground pipelines only.

The United Kingdom has developed 3 classifications, with the separation distances to existing occupied structures being a function of the specific classification, the operating pressure and the outside diameter of the pipe. However, the minimum separation distance requirements decrease with increased population density (which allows for siting of pipelines in heavily populated areas).

Since 1972, the Health and Safety Executive (HSE) of Great Britain has been providing advice to planning authorities concerning the potential risks and effects of hazards to neighboring populations. However, HSE guidance for construction near pipelines is still under development (6).

As described in information provided by Hans van Poelje of N.V. Nederlandse Gasunie (personal communication, November 30, 1995), the Dutch have established four area classes, and have identified two separation distances - effect and building distances. The Dutch zoning guidelines recommend building outside of the effect distance if possible, and it is the area inside the effect distance that is actually classified from 1 (low population density) to 4 (high density). The effect and building distances depend on the diameter and pressure of the pipe, and are based on calculated contours of individual risk
(which vary from $10^{-8}$ to $10^{-6}$ ). Individual risk is defined by the Dutch (H. van Poelje: personal communication, September 30, 1996) as being the probability per year of the death of a person who stands 24 hours a day undressed at a certain distance from the pipeline.

In general, both the United States and foreign countries address the establishment of separation distances either directly, or indirectly through the designation of various location or population classes. In those instances where a defined allowable ROW has been established, the width of the ROW from the pipe centerline is relatively small (less than 100 feet) and is generally intended to protect the pipe rather than the public. Separation distances have also been developed by some countries based on the concept of risk assessment. This dissertation describes the development of a methodology to estimate actual separation distances required, based upon the effects of a rupture of a natural gas transmission pipeline of a specified diameter and operating pressure. The methodology is developed, in part, with data collected from prior investigations of actual natural gas transmission pipeline accidents.

The dissertation is presented in seven chapters. Chapter 1 is an introduction. Chapter 2 discusses the literature review that has been performed; Chapter 3 provides the objective of the dissertation. Chapter 4 develops a new methodology by which safe separation distances can be determined. This technique is then compared, in Chapter 5, to actual pipeline accident data, and to the work of previous researchers. The results of the dissertation are presented and discussed in Chapter 6, while Chapter 7 provides conclusions and recommendations for future work.

## CHAPTER 2

## LITERATURE REVIEW

The purpose of the literature review was to gain an understanding of prior research concerning both natural gas pipeline accidents and the determination of safe separation distances from natural gas transmission pipelines. The literature review consisted of the following activities:

- Review of NTSB files in Washington, DC;
- Review of Commodity Pipeline Occurrence Reports, prepared by the Canadian TSB;
- Review of databases compiled by the New Jersey Institute of Technology's (NJIT's) Institute for Transportation containing relevant citations of pipeline accidents and safe separation distances;
- Correspondence with the European Gas Pipeline Incident Data Group located in the Netherlands;
- Review of various texts pertaining to pipelines, hazard assessment, heat transfer and fluid flow;
- Review of technical journals of the oil and gas pipeline industry; and
- Discussions with the NTSB, the Canadian TSB, the Canadian National Energy Board, the United States Environmental Protection Agency, the American Gas Association, the Research and Special Programs Administration, the Gas Research Institute, and the Township of Edison regarding pipeline accidents and safe separation distances.

In general, the literature review revealed that there has been limited studies to date concerning the determination of safe separation distances. Before discussing the prior research, however, the properties of natural gas and the principles of fluid flow are discussed below.

### 2.1 Natural Gas and its Properties

Natural gas is composed primarily of methane. A typical commercial field gas has the following molar composition (7):

Table 2.1
Composition of Natural Gas

| Component | Mole <br> Percent |  |
| :--- | :--- | :--- |
| Oxygen | 0.0008 <br> Nitrogen | 1.498 |
| Carbon Dioxide | 1.073 |  |
| Methane | 83.266 |  |
| Ethane | 9.608 |  |
| Propane | 3.597 |  |
| Isobutane | 0.3414 |  |
| Normal Butane | 0.4581 |  |
| Isopentane | 0.0403 |  |
| Normal Pentane | 0.0342 |  |
| Hexane | 0.0046 |  |
| Heptane | 0.0003 |  |
| Octane | 0.0001 |  |
| Toluene | 0.0002 |  |

Natural gas generally has a gross heat content of $1,000 \mathrm{Btu} / \mathrm{ft}^{3}$ measured dry at $60^{\circ} \mathrm{F}$ and 30 inches (or one atmosphere) of mercury (8). Since $60^{\circ} \mathrm{F}$ and 1 atmosphere are considered to be respectively the standard conditions for temperature and pressure in the gas industry (NTSB: personal communication, February 23, 1996), the heat content can be expressed as 1,000 Btu per standard cubic foot (scf).

Since methane is the primary constituent of natural gas, the gas mixture will exhibit physical properties similar to the properties for pure methane. At I atmosphere, methane has a freezing point in air of $-296.45^{\circ} \mathrm{F}$ and a boiling point of $-258.68^{\circ} \mathrm{F}(8)$. Therefore, at standard conditions, methane exists as a gas. Thermodynamically, the ratio of heat capacity at constant pressure $(\mathrm{Cp})$ to heat capacity at constant volume $(\mathrm{Cv})$ for methane varies from 1.39 at $-150^{\circ} \mathrm{F}$ to 1.25 at $150^{\circ} \mathrm{F}(8)$.

Methane is a combustible gas, and has a lower flammability limit of approximately $5 \%$ by volume (8). The upper flammability limit is approximately $15 \%$ by volume (8). This implies that mixtures of methane and air will burn if the mixture contains between $5 \%$ and $15 \%$ (by volume) of methane. Outside of this range, the mixture is either too lean or too rich to burn.

Natural gas containing hydrogen sulfide is called sour gas. An example of a sour gas mixture (9) is presented in Table 2.2:

Table 2.2
Composition of Sour Gas

| Component | Mole <br> Percent |
| :--- | :--- |
| Nitrogen | 2.3046 |
| Carbon Dioxide | 1.1616 |
| Hydrogen Sulfide | 9.0791 |
| Methane | 58.7939 |
| Ethane | 12.7407 |
| Propane | 10.3420 |
| Isobutane | 1.1425 |
| Normal Butane | 3.0209 |
| Isopentane | 0.6091 |
| Normal Pentane | 0.5910 |
| Hexane | 0.1217 |
| Other Components | 0.0929 |

Gases (including natural gas) are considered compressible fluids. Therefore, the following discussion will present the mathematical relationships that describe compressible fluid flow.

### 2.2 Compressible Fluid Flow

### 2.2.1 Choked Conditions

When a vessel such as a tank is accidentally punctured (resulting in a small hole) or intentionally blown down (through a nozzle), the gas (provided that it is at a sufficient pressure) may leave the vessel at its sonic velocity. The sonic velocity is the speed of sound in the gas, and is the maximum possible speed through the puncture or nozzle. A
"choked" condition is established if the gas leaves the vessel at a velocity equal to the sonic velocity. In order for choked flow to occur, the following criterion (assuming an ideal gas) must be satisfied (10):

$$
\begin{equation*}
\left(\mathrm{p}_{2} / \mathrm{p}_{1}\right)_{\mathrm{c}}=(2 /(\mathrm{k}+1))^{\mathrm{k} /(\mathrm{k}-1)} \tag{2-3}
\end{equation*}
$$

where:
$\mathrm{p}_{1}=$ Absolute vessel pressure (slug/ft s${ }^{2}$ )
$\mathrm{p}_{2}=$ Absolute pressure in the minimum section (slug/ft s${ }^{2}$ )
$\mathrm{k}=$ Gas specific heat ratio, $\mathrm{Cp} / \mathrm{Cv}$
$c=$ Critical conditions
If the pressure ratio $p_{2} / p_{1}$ (absolute pressure outside of the vessel to inside the vessel) is greater than the critical ratio and isentropic conditions are assumed (i.e., frictionless conditions where there is no exchange of heat), then the mass flow rate out of the vessel can be calculated as follows (10):

$$
\begin{equation*}
\mathrm{m}^{*}=\mathrm{A}_{2}\left\{(2 \mathrm{k} /(\mathrm{k}-1)) \mathrm{p}_{1}\left(\rho_{1}\right)\left(\left(\mathrm{p}_{2} / \mathrm{p}_{1}\right)^{2 / \mathrm{k}}-\left(\mathrm{p}_{2} / \mathrm{p}_{1}\right)^{(\mathrm{k}+1) / k)}\right\}^{1 / 2}\right. \tag{2-4}
\end{equation*}
$$

where:
$\mathrm{A}_{2}=$ Area of minimum section of the puncture or nozzle $\left(\mathrm{ft}^{2}\right)$
$\rho_{1_{*}}=$ Density of gas inside vessel (slug/ft ${ }^{3}$ )
$\mathrm{m}^{*}=$ Mass flow rate (slug/s)
If the pressure ratio $p_{2} / p_{1}$ is less than the critical ratio, then the mass flow rate is found from (16):

$$
\begin{equation*}
\mathrm{m}^{*}=\left(\mathrm{A}_{2} \mathrm{p}_{1} /\left(\mathrm{T}_{1}\right)^{1 / 2}\right)\left\{(\mathrm{k} / \mathrm{R})(2 / \mathrm{k}+1)^{(\mathrm{k}+1) /(\mathrm{k}-1)}\right\}^{1 / 2} \tag{2-5}
\end{equation*}
$$

where
$\mathrm{R}=\operatorname{Gas}$ constant ( $5.98 \mathrm{ft} \mathrm{lb} /$ slug $^{\circ} \mathrm{R}$ )
$\mathrm{T}_{1}=$ Gas temperature $\left({ }^{\circ} \mathrm{R}\right)$

### 2.2.2 Isothermal Pipe Flow

Gas flow in pipelines can be isothermal. For isothermal conditions to exist (i.e., the gas temperature is constant), the heat transferred from the gas through the pipe walls is balanced by the heat generated by friction (10). Such conditions can occur in uninsulated pipes where velocities are much less than sonic and where the temperatures inside and outside the pipeline are of the same order.

The following equation relates upstream and downstream pressures for isothermal
flow (16):

$$
\begin{equation*}
\mathrm{p}_{1}^{2}-\mathrm{p}_{2}^{2}=\left(\mathrm{m}^{* 2} \mathrm{RT} / \mathrm{A}^{2}\right)\left(2 \ln \left(\mathrm{p}_{1} / \mathrm{p}_{2}\right)+\mathrm{f}(\mathrm{l} / \mathrm{d})\right) \tag{2-6}
\end{equation*}
$$

where:
$\mathrm{p}_{1}=$ Upstream pressure (slug/ft s${ }^{2}$ )
$\mathrm{p}_{2}=$ Downstream pressure, (slug/ft s${ }^{2}$ )
$\mathrm{A}=\mathrm{Cross}$-sectional area of pipeline $\left(\mathrm{ft}^{2}\right)$
$\mathrm{T}=$ Gas temperature $\left({ }^{\mathrm{O}} \mathrm{R}\right)$
$\mathrm{R}_{*}=$ Gas constant ( $5.98 \mathrm{ft} \mathrm{lb/slug}{ }^{\mathrm{o}} \mathrm{R}$ )
$\mathrm{m}=$ Mass flow rate (slug/s)
$\mathrm{f}=$ Moody friction factor
$1=$ Distance between upstream and downstream locations (ft)
$\mathrm{d}=$ Pipeline diameter (ft)
The above equation can also be written in terms of $k$ (heat capacity ratio) and the upstream and downstream Mach numbers $\left(M_{1}, M_{2}\right.$ respectively, where $M$ is the ratio of the gas velocity to the sonic velocity) (10):

$$
\begin{equation*}
\left(M_{1}^{2} / M_{2}^{2}\right)=1-\mathrm{kM}_{1}^{2}\left(2 \ln \left(\mathrm{M}_{2} / \mathrm{M}_{1}\right)+\mathrm{f}(1 / \mathrm{d})\right) \tag{2-7}
\end{equation*}
$$

These equations are applicable only for the range

$$
\begin{equation*}
M_{1}<M_{2} \leq(1 / k)^{1 / 2} \tag{2-8}
\end{equation*}
$$

### 2.2.3 Adiabatic Pipe Flow

As described by Vennard \& Street (10), adiabatic conditions (i.e., conditions in which there is no heat transfer) can occur in pipes that are either insulated or of short length. Under these conditions, the Mach number increases downstream in subsonic flow, but decreases downstream in supersonic flow. In both subsonic and supersonic flow, the Mach number tends toward unity, with $M=1$ being the limiting condition. Given an upstream condition, the distance "I" downstream of that location can be increased until $M=1$ at the discharge location. Any further increases in " 1 " cannot be made without altering the upstream conditions or creating a shock wave. The relevant equation used for performing calculations is:

$$
\begin{equation*}
\left(d \mathrm{M}^{2} / \mathrm{M}^{2}\right)=\mathrm{k} \mathrm{M}^{2}\left\{\left(1+((\mathrm{k}-1) / 2) \mathrm{M}^{2}\right) /\left(1-\mathrm{M}^{2}\right)\right\} \mathrm{f}(d \mathrm{l} / \mathrm{d}) \tag{2-9}
\end{equation*}
$$

With this understanding of natural gas properties and compressible flow, the work of several investigators regarding pipeline blowdown can be discussed.

### 2.3 Prior Research Pertaining to Pipeline Blowdown

### 2.3.1 Wilson

In his 1979 study for Alberta Environment, Wilson (11) developed a method for evaluating ground-level transient gas releases for the purpose of assessing risks from sour gas pipeline ruptures. As described by Wilson, automatic block valves are used to isolate the ruptured pipeline segment. Furthermore, during blowdown, the following time periods exist:

- Soon after the rupture, an expansion wave moves down the pipeline at a velocity equal to the speed of sound. The exiting mass flow can be thought of as an emission from a source of gas at the initial temperature and pressure of the pipeline. Additionally, the gas is considered to flow through a pipe length which increases with time at a rate equal to the speed of sound.

Once the wave reaches the end of the pipe, the system will behave as a gas source of constant size which exits through a hole at one end.

Wilson performs both adiabatic and isothermal analyses of the initial moving wave transient. Furthermore, the author also performs an analysis for an isothermal release during the subsequent condition of constant length. An assumption is made that the system is quasi-steady state, which means that fluid acceleration forces are small as compared to pressure forces. The author concludes that the mass release rate models developed through these analyses are difficult to use, and recommends using the following double exponential model developed by R. P. Bell:

$$
\begin{equation*}
\mathrm{m}^{*} \mathrm{e}=\left(\mathrm{m}^{*} \mathrm{oe}^{\prime} /(1+\alpha)\right)\left[\exp \left(-\mathrm{t} /\left(\alpha^{2} \beta\right)\right)+\alpha \exp (-\mathrm{t} / \beta)\right] \tag{2-10}
\end{equation*}
$$

where, in dimensions of mass $(M)$, length $(\mathrm{L})$ and time $(\mathrm{t})$, the terms are:
$\mathrm{m}^{*} \mathrm{e}=$ Mass release rate of gas at time $\mathrm{t}(\mathrm{M} / \mathrm{t})$
$\mathrm{m}^{*} \mathrm{oe}=$ Initial mass release rate calculated using Eq. 2-5 (M/t)
$\alpha=W_{o} /\left[\beta \mathrm{m}^{*}{ }_{\mathrm{oe}}\right]$ (dimensionless)
$W_{0}=$ Total mass of gas in pipeline $(M)$
$\beta=$ Final release time constant $(t)$

In his later work (12), Wilson describes the initial mass release rate as a fraction of the value obtained by Equation 2-5:

$$
\begin{equation*}
\mathrm{m}_{\mathrm{oe}}^{*}=\mathrm{K}\left(\mathrm{~m}_{\mathrm{oe}}^{*}\right. \text { found using Eq. 2-5) } \tag{2-11}
\end{equation*}
$$

Where $K$ = equals a constant less than I, which accounts for the resistance to flow caused by the overlying soil resting on the pipe. Wilson compares the results from two full scale pipe rupture tests, and concludes that a $K$ value of 0.5 provides the best fit of the double exponential model.

In a later work concerning the risk of exposure concentrations in fluctuating plumes (13), Wilson states that "predicting the behavior of a toxic gas plume from a pipeline rupture is one of the most complex and difficult problems in atmospheric diffusion." He cites complicating factors including the high pipeline pressures which make the exiting gas supersonic and which results in large changes in gas temperature and density. Furthermore, he indicates that the plume is sensitive to changes in terrain, downwind obstacles and atmospheric stability. He also mentions that a "saving grace in this messy problem" is that the release will be of short duration ( 30 minutes or less).

### 2.3.2 Groves, Bishnoi and Wallbridge

The authors (7) discuss the calculation of wave velocity in a ruptured pipeline. As previously mentioned, the rupture of a high pressure gas pipeline results in the creation of an expansion (or decompression) wave. This wave moves away from the rupture as the
gas flows toward it. Since natural gas is actually a mixture of condensible gases, decompression can result in condensation of individual compounds. Several of the assumptions that the authors make include the following:

- The flow is isentropic and nonviscous;
- The rupture is instantaneous and over the full cross-section of the pipe;
- Condensation occurs at constant temperature and pressure;
- After condensation, the liquid mixture is in equilibrium with the gas mixture; and - The gas mixture is described by the Soave-Redlich-Kwong equation of state, which is referenced in the authors' document.

The authors present an iterative procedure to find the wave velocity, and compare theoretical predictions to experimental data for pure methane, argon and a natural gas mixture. The authors conclude by stating that it is possible to calculate the decompression wave behavior from pipeline ruptures containing pure gases or gas mixtures.

### 2.3.3 Fannelop and Ryhming

The work presented by these authors (14) involves the sudden and complete break of a high-capacity gas transmission line. The case which is used is an underwater pipeline 90 miles long, 36 inches in diameter, with inlet and outlet pressures of 133 atm and 55 atm , respectively. The ambient pressure outside the pipeline is 6 atm , and the line contains

7,000 tons of gas. Isothermal flow of a perfect (i.e., ideal) gas is assumed, with the flow also being one-dimensional, unsteady, compressible and with friction. The following equations are then presented:

Continuity:

$$
\begin{equation*}
\partial \rho / \partial t+(\partial / \partial x) \rho u=0 \tag{2-12}
\end{equation*}
$$

Momentum:

$$
\begin{equation*}
\partial u / \partial t+u(\partial u / \partial x)+(1 / \rho)(\partial \rho / \partial x)+2 f u^{2} / D=0 \tag{2-13}
\end{equation*}
$$

Equation of state:

$$
\begin{equation*}
p=\rho R T_{0} \tag{2-14}
\end{equation*}
$$

where, in dimensions of mass (M), length (L), time ( $t$ ) and temperature ( T ), the various terms are:
$\rho=$ gas density $\left(\mathrm{M} / \mathrm{L}^{3}\right){ }^{3}$
$\mathrm{R}=$ gas constant $\left(\mathrm{L}_{2}^{2} / \mathrm{Tt}^{2}\right)$
$\mathrm{p}=$ pressure $\left(\mathrm{M} / \mathrm{Lt}^{2}\right)$
$t=$ time $(t)$
$\mathrm{u}=$ gas velocity $(\mathrm{L} / \mathrm{t})$
$x=$ length coordinate $(L)$
$f=$ friction factor $(f=0.002)$
$D=$ pipeline diameter ( L )
$\mathrm{T}_{0}=$ seawater temperature $(\mathrm{T})$
The authors describe different time regimes following the rupture, which require different methods of solution. The "early time" regime begins with an inviscid regime followed by a viscous expansion process. This carly time regime lasts from the time of rupture until approximately 25 seconds or the waves penetrate 10,000 meters of pipeline.

The "intermediate-time" regime lasts from approximately 25 seconds until the pressure
peak has moved to the closed end of the pipeline, at which time the "late-time" regime begins. When considering the low pressure end of the rupture, an intermediate region is not considered (since the pressure decreases monotonically toward the open end).

The authors acknowledge that a general solution of the above nonlinear partial differential equations is too complicated for routine engineering applications. They present an integral method which can be used to find the solution for the late time regime (which represents most of the outflow from the broken end). The authors also indicate that this method can be used (with modification) for the intermediate regime.

### 2.3.4 Flatt

In his 1985 work, Flatt (15) provided a second-order algorithm which is applicable to the inverse marching method of characteristics. This algorithm can be used to solve unsteady compressible one-dimensional flows in which the momentum equation has a high $f(\mathrm{~L} / \mathrm{D})$ value and where the flow is in the high subsonic range. The author indicates that the method involves only one iteration.

Flatt acknowledges in another study (16) that the rupture problem is highly non-linear, making the finding of analytical solutions unlikely. He applies the inverse marching method of characteristics to the pipeline rupture presented in (14). The author assumes that initial conditions are of isothermal flow, while the unsteady flow generated after the rupture is adiabatic and one-dimensional. Additionally, the gas is considered perfect.

### 2.3.5 Lang and Fannelop

The pipeline rupture described in previously noted studies $(14,16)$ is again analyzed in a subsequent study (17). The partial differential equations are reduced by the authors to a set of ordinary differential equations by procedures used in the Method of Weighted Residuals. The reduced equations are then integrated using various numerical methods (Finite Element, Spectral Galerkin and Spectral Collocation). The authors assume the flow to be one-dimensional, compressible, viscous and isothermal, with the conditions of primary interest being those in the pipeline segment associated with the high pressure end. The authors state that the best results in terms of stability, accuracy and computing time are obtained using the Spectral Collocation Method.

### 2.3.6 Ryhming

The author considers the early time regime associated with a pipeline rupture (18). Specifically, the process of the expansion wave imparting motion to the gas particles against wall friction is studied. The governing equations are those provided in (14), and are solved using a matched asymptotic expansion procedure. The author also assumes isothermal flow conditions and a perfect gas.

While the author acknowledges that numerically correct solutions are obtained by considering the full adiabatic process, the isothermal assumption results in simpler equations and very little is lost in terms of the gas acceleration process.

### 2.3.7 Picard and Bishnoi

The authors use a real-fluid nonisentropic decompression model to evaluate the rupture of a sour gas pipeline (9). This model is used for two-phase (gas-liquid) multi component mixtures, and is based on unsteady nonisentropic one-dimensional compressible fluid flow. The model consists of characteristic and compatibility equations that are solved numerically using an explicit inverse marching method. Real fluid behavior is modeled using the Peng and Robinson equation of state.

The pipeline considered is part of a gas gathering system. It is uninsulated, 168.3 mm diameter and buried 42 cm below the ground (ground temperature is 275 K ). The pre-rupture temperature of the fluid is 303 K , while the corresponding pressure is 11,000 kPa . When the pipeline is ruptured, it is assumed that an upstream emergency valve closes, restricting the blowdown length to $1,000 \mathrm{~m}$.

The authors compare the results with those obtained using the assumption of a perfect gas, and conclude that a perfect gas assumption will underestimate the release rate by 30 to $45 \%$ while the total amount of fluid released may be underestimated by $50 \%$. Additionally, the authors point out that perfect gas theory does not consider the effects of condensation.

### 2.3.8 Olorunmaiye and Imide

The authors (19) model the flow in the underwater pipeline that has been described previously. They assume unsteady, one-dimensional isothermal flow. The non-dimensional forms of the partial differential equations provided above are solved using the method of characteristics. The authors indicate that the flow rate predicted at the broken end of the pipe is $18 \%$ lower than that predicted by adiabatic flow theory, but that the flow rate agrees with the results of Lang and Fannelop.

### 2.3.9 Other Work

The Pipe Line Rules of Thumb Handbook (8) provides the following equation to use when calculating the volume of gas lost through a puncture or blowdown:

$$
\begin{equation*}
\mathrm{Q}=\mathrm{D}^{2} \mathrm{P}_{1} \tag{2-15}
\end{equation*}
$$

where:
$\mathrm{Q}=$ Volume of gas in $\mathrm{Mcf} / \mathrm{hr}$ at a pressure of $14.9 \mathrm{psi}, 60^{\circ} \mathrm{F}$ with a specific gravity of 0.60
$D=$ Diameter of the nipple or orifice in inches
$P_{1}=$ Absolute pressure in psi at some nearby point upstream from the opening

### 2.4 Prior Research Pertaining to Turbulent Jet Flames and Flame Geometry

### 2.4.1 Turbulent Jet Flames

As described by the National Fire Protection Association and the Society of Fire Protection Engineers (20), turbulent jet flames can form as the result of either an accidental release of hydrocarbon vapors or the intentional burning of waste gases in a flare. The flame is called turbulent based on the Reynolds number of the gas jet. A jet
fuel is assumed to exit a nozzle that is pointed in an upward position. The fuel is discharged into a surrounding medium that contains an oxidant. Starting at a low flow rate, the flame height increases with increasing flow rate. As the flow is increased, a maximum height is reached and then shortens. However, before reaching the maximum height, the flame begins to flicker at the top. This point changes the flame from a laminar diffusion flame to a transition flame. As the flow rate is increased further, the flickering spreads down and stops several diameters from the burner. When this happens, the flame is said to be turbulent, and the flame height becomes independent of flow rate.

### 2.4.2 Flame Geometry

Two important flame characteristics are the flame height and the amount of flame tilt in a cross-wind. Hawthorne, Weddell and Hottel (21) have developed the following relationship to predict the height of a turbulent flame in still air:

$$
\begin{equation*}
\mathrm{L} / \mathrm{D}=\left(5.3 / \mathrm{C}_{\mathrm{T}}\right)\left\{\left(\mathrm{T}_{\mathrm{F}} /\left(\alpha_{\mathrm{T}} \mathrm{~T}_{\mathrm{N}}\right)\right)\left(\mathrm{C}_{\mathrm{T}}+\left(1-\mathrm{C}_{\mathrm{T}}\right)\left(\mathrm{M}_{\mathrm{S}} / \mathrm{M}_{\mathrm{N}}\right)\right\}^{1 / 2}\right. \tag{2-16}
\end{equation*}
$$

where, in dimensions of mass (M), length (L), time ( $t$ ) and temperature ( $T$ ), the various terms are:
$\mathrm{L}=$ visible flame length (L)
$\mathrm{D}=$ nozzle diameter ( L )
$\mathrm{T}_{\mathrm{F}}=$ adiabatic flame temperature (absolute T )
$\mathrm{T}_{\mathrm{N}}=$ temperature of nozzle fluid (absolute T )
$\mathrm{M}_{\mathrm{S}}=$ molecular weight of surrounding fluid (dimensionless)
$\mathrm{M}_{\mathrm{N}}=$ molecular weight of nozzle fluid (dimensionless)
$\mathrm{C}_{\mathrm{T}}=$ mol fraction of nozzle fluid in the unreacted stoichiometric mixture (decimal)
$\alpha_{T}=$ mols of reactants/mols products for the stoichiometric mixture (decimal)
facing 22
Note: Illustration Based on Information Provided in References Numbered 20 and 22 X (Direction of Cross-wind)
Figure 2.1 - Flame Shape with a Cross-wind

The work of Brzustowski, as presented in (20) and Gollahalli, et al (22) provides a procedure to determine flame shape with a cross-wind (see Figure 2.1). The procedure begins by calculating the dimensionless lean limit concentration:

$$
\begin{equation*}
\overline{\mathrm{C}}_{\mathrm{L}}=\mathrm{C}_{\mathrm{L}}\left(\mathrm{u}_{\mathrm{j}} / \mathrm{u}_{\mathrm{w}}\right)\left(\mathrm{M}_{\mathrm{f}} / \mathrm{M}_{\mathrm{a}}\right) \tag{2-17}
\end{equation*}
$$

where:
$\overline{\mathrm{C}}_{\mathrm{L}}=$ lean limit concentration (expressed as a decimal)
$\mathrm{u}_{\mathrm{j}}=$ velocity of jet
$u_{w}=$ velocity of wind
$\mathrm{M}_{\mathrm{f}}=$ molecular weight of fuel
$M_{a}=$ molecular weight of air
If $\bar{C}_{L} \leq 0.5$, the curvilinear distance $\bar{S}_{L}$ is determined from:

$$
\begin{equation*}
\overline{\mathrm{S}}_{\mathrm{L}}=2.04\left(\overline{\mathrm{C}}_{\mathrm{L}}\right)^{-1.03} \tag{2-18}
\end{equation*}
$$

If $\overline{\mathrm{C}}_{\mathrm{L}}>0.5$, then:

$$
\begin{equation*}
\left.\overline{\mathrm{S}}_{\mathrm{L}}=2.71 \overline{\mathrm{C}}_{\mathrm{L}}\right)^{-0.625} \tag{2-19}
\end{equation*}
$$

If $\bar{S}_{L}>2.35$, then the dimensionless coordinate $\bar{X}_{L}$ is found from:

$$
\begin{equation*}
\bar{X}_{L}=\bar{S}_{\mathrm{L}}-1.65 \tag{2-20}
\end{equation*}
$$

If $\overline{\mathrm{S}}_{\mathrm{L}}<2.35$, then $\overline{\mathrm{X}}_{\mathrm{L}}$ is determined from:

$$
\begin{equation*}
\overline{\mathrm{S}}_{\mathrm{L}}=1.04 \overline{\mathrm{X}}_{\mathrm{L}}^{2}+2.05 \overline{\mathrm{X}}_{\mathrm{L}}^{0.28} \tag{2-21}
\end{equation*}
$$

The dimensionless rise $\bar{Z}_{L}$ is calculated from:

$$
\begin{equation*}
\overline{\mathrm{Z}}_{\mathrm{L}}=2.05\left(\overline{\mathrm{X}}_{\mathrm{L}}\right)^{0.28} \tag{2-22}
\end{equation*}
$$

The dimensional coordinates of the flame tip are then found from the following:

$$
\begin{align*}
& X=\bar{X}_{L} d_{j}\left(\rho_{j} / \rho_{a}\right)^{1 / 2}\left(u_{j} / u_{w}\right)  \tag{2-23}\\
& Z=\bar{Z}_{L} d_{j}\left(\rho_{j} / \rho_{a}\right)^{1 / 2}\left(u_{j} / u_{w}\right) \tag{2-24}
\end{align*}
$$

facing 23

Note: Illustration Based on Information Provided in References Numbered 20 and 23
Figure 2.2 - Jet Flame Shape Parameters
where:
$\mathrm{d}_{\mathrm{j}}=$ Jet diameter
$\rho_{j}=$ Density of jet
$\rho_{a}=$ Density of air
The work of Kalghatgi presented in (20) and (23) provides a means by which the shape of a jet flame in a cross-wind can be described. He described the flame by a frustum of a cone (see Figure 2.2) with the following parameters:
$\mathrm{W}_{1}=$ flame base width
$\mathrm{W}_{2}=$ flame tip width
$\mathrm{L}_{\mathrm{BV}}=$ vertical length of the flame tip from the plane of the burner
$\alpha=$ angle subtended by the flame with respect to the vertical
$\alpha_{B}=$ angle subtended by the burner tip and the tip of the flame with respect to the vertical

Kalghatgi also defined the burner source diameter as being

$$
\begin{equation*}
D_{S}=D\left(\rho_{j} / \rho_{\mathrm{a}}\right)^{1 / 2} \tag{2-25}
\end{equation*}
$$

where $D$ is the actual burner diameter. A velocity ratio, $R$, is also defined as:

$$
\begin{equation*}
\mathrm{R}=\text { cross-wind speed }(\mathrm{U}) \text { jet velocity }\left(\mathrm{U}_{\mathrm{j}}\right) \tag{2-26}
\end{equation*}
$$

The following equations are then used to find the geometric parameters (for the range $0.02<R<0.25$ ):

$$
\begin{gather*}
\alpha_{B}=94-(1.6 / R)-35 R \text { (degrees) }  \tag{2-27}\\
\alpha=94-(1.1 / R)-30 R(\text { degrees })  \tag{2-28}\\
\left(L_{B V} / D_{S}\right)=6+(2.35 / R)+20 R  \tag{2-29}\\
\left(W_{2} / D_{S}\right)=80-(0.57 / R)-570 R+1470 R^{2}  \tag{2-30}\\
\left(W_{1} / D_{S}\right)=49-(0.22 / R)-380 R+950 R^{2} \tag{2-31}
\end{gather*}
$$

facing 24

Figure 2.3 - Conceptual Illustration of the Point Source Model

### 2.4.3 Heat Transfer Models

The major source of damage or injury from large open hydrocarbon fires is from thermal radiation (20). The following discussion will present several models that can be used to predict the thermal radiation from jet flames.
2.4.3.1 Point Source Model: The point source model assumes that the source of radiant heat can be modeled by positioning a point source at the center of the flame (see Figure 2.3). The equation which is used to calculate the radiant heat flux is (20):

$$
\begin{equation*}
\mathrm{q}^{\prime \prime}=\mathrm{Fq}_{\mathrm{rel}} \cos \theta /\left(4 \pi \mathrm{D}^{2}\right) \tag{2-32}
\end{equation*}
$$

where:
$\mathrm{q}^{\prime \prime}=$ incident radiant flux, $\mathrm{kW} / \mathrm{m}^{2}$ (or Btw/hr $\mathrm{ft}^{2}$ )
$\mathrm{D}=$ distance from flame center to observer, m (or ft )
$q_{\text {rel }}=$ energy release rate, kW (or Btu/hr)
$\mathrm{F}=$ fraction of combustion energy resulting in radiation
$\theta=$ angle between normal to the surface receiving the thermal radiation and line of sight from flame center

Oenbring and Sifferman (24) provide the following form of the point source
model to use when calculating radiant heat flux from a flare:

$$
\begin{equation*}
\mathrm{K}=\mathrm{FQ} /\left(4 \pi \mathrm{D}^{2}\right) \tag{2-33}
\end{equation*}
$$

where:
$\mathrm{K}=$ radiant heat flux from a flame $\left(\mathrm{Btu} / \mathrm{hr}_{\mathrm{ft}}{ }^{2}\right)$
$F=$ fraction of total heat radiated
$Q=$ total heat content of flared gas (Btu/hr)
It can be seen from both forms of the model that the heat flux decreases with the square of the distance (i.e., there is an inverse square relationship).
facing 25

2.4.3.2 Line Source Model: The NFPA and SFPE (20) explain that in the line source model, an elemental length of flame is assumed to radiate similar to a point source. The total radiant heat flux is calculated by integrating the flux from an elemental source over the length of the flame.

In addition, Fumarola, et al (25) have developed the following equation for use when the thermal radiation is evenly distributed along the flame axis:

$$
\begin{equation*}
\mathrm{q}=0.043 \mathrm{fQ} \int\left[\mathrm{Ax} 0^{0.36}+\mathrm{h}-\overline{\mathrm{z}}\right] /\left[\mathrm{x}\left\{\left(\mathrm{Ax}{ }^{0.36}+\mathrm{h}-\overline{\mathrm{z}}\right)^{2}+(\mathrm{x}-\overline{\mathrm{x}})^{2}\right\}^{3 / 2}\right] \mathrm{dx} \tag{2-34}
\end{equation*}
$$

In Equation 2-34, $A$ is equal to $3.1\left(d_{j} R\right)^{0.64}$, R is equal to $\left(\mathrm{u}_{\mathrm{j}} / \mathrm{u}_{\mathrm{a}}\right)\left(\rho_{j} / \rho_{\mathrm{a}}\right)^{1 / 2}$, the limits of integration are from $x$ equals zero to $x$ equals $x_{t}$, and the following nomenclature is used:
$\mathrm{d}_{\mathrm{j}}=$ diameter of flare stack, m
$h=$ height of flare stack, $m$
$f=$ fraction of radiant heat release
$\mathrm{Q}=$ total heat release, $\mathrm{kcal} / \mathrm{s}$
$\mathrm{q}=$ thermal radiation at receptor, $\mathrm{kcal} / \mathrm{s} \mathrm{m}^{2}$
$x=$ downstream distance, $m$
$z=$ cross-stream distance, $m$
$\bar{x}=x$-coordinate of receptor, $m$
$\bar{z}=z$-coordinate of receptor, $m$
$\mathrm{u}=$ velocity, $\mathrm{m} / \mathrm{s}$
$a=$ ambient air condition
$j=$ discharge condition
$t=$ end of the flame axis
A conceptual illustration of the line source model is provided in Figure 2.4.

### 2.4.3.3 Tilted Cylinder Model: While the point source model can be used to determine

 the thermal radiation from large turbulent flames, the model breaks down at locations close to the flame (20). Another model that can be used is the tilted cylinder model (20). This model approximates the flame by a cylindrical radiating surface. The radiant heat flux can be calculated from the following equation:$$
\begin{equation*}
\mathrm{q}=\mathrm{E} \tau \mathrm{~F} \tag{2-35}
\end{equation*}
$$

where:
$E=$ surface emissive power of the flame, $\mathrm{W} / \mathrm{m}^{2}$ (or $\left.\mathrm{Btu} / \mathrm{hr} \mathrm{ft}^{2}\right)$
$\mathrm{F}=$ geometric view factor
$\tau=$ atmospheric transmissivity
The surface emissive power of the flame can be found from the equation:

$$
\begin{equation*}
E=E_{b b}\left(1-e^{-k L}\right) \tag{2-36}
\end{equation*}
$$

where:
$\mathrm{E}_{\mathrm{bb}}=$ equivalent blackbody emissive power (same units as E )
$\mathrm{k}=$ extinction coefficient, $1 / \mathrm{m}$ (or $1 / \mathrm{ft}$ )
$\mathrm{L}=$ effective path length, m (or ft )

### 2.5 Prior Research in the Field of Hazard Analysis

### 2.5.1 Typical Models

There are numerous computer models that have been developed for use in the field of hazard analysis. These models are generally used to predict the fate and transport of hazardous materials that have been accidently released into the environment. Several of these models are described below.

ALOHA (Areal Locations of Hazardous Atmospheres) (26) was developed by the National Oceanic and Atmospheric Administration and the United States Environmental Protection Agency. The software is intended for use by emergency responders. ALOHA performs dispersion modeling for neutrally buoyant and heavy gases and (for a variety of chemicals) can model different sources such as tanks and pipes.

FLACS (Flame Acceleration Simulator) (27) has been used to predict gas explosions in complex geometries. The computer model can be used in the design of process areas (including offshore platforms). FLACS calculates parameters such as explosion pressure and flow characteristics as a function of time and space for different geometries and explosion situations.

SAFEMODE (Safety Assessment for Effective Management Of Dangerous Events) and Micro HACS (Hazard Assessment Computer System) were developed for the United States Coast Guard (28). This system can model releases from sources such as pressurized storage tanks, bulk storage tanks, barge, rail car and road transports) and predict fate and transport (such as explosion effects and pool fire thermal radiation) for over 1200 chemicals.

The Hazard model can be used with pipeline or tank storage failures (29). Predictions can be made of the specific locations of hazardous concentration limits (whether they are flammable or toxic limits) in the event of a release. The model treats gas pipeline failures using one-dimensional gas dynamics.

WHAZAN (World Bank Hazard Analysis) was developed by Technica International Ltd. in collaboration with the World Bank (30). The overall program is comprised of a series of consequence models. The models can predict the outflow of a chemical, its behavior immediately after release, its dispersion in the atmosphere, and fire or explosion events.

### 2.5.2 Research Performed by the European Gas Pipeline Incident Data Group

In Europe, there are 8 major natural gas transportation companies that are collectively called the European Gas Pipeline Incident Data Group (EGIG). In the 1960's and early 1970's, European governments based their regulatory actions on accident data from the United States. The European gas companies believed it was more appropriate to have their regulations based on European data, and so they decided to establish their own database. EGIG collects information on the technical conditions of the incident, such as the cause of the incident, incident frequencies, the diameter of the pipeline, the pipe wall thickness, and the thickness of the pipeline cover. The type of information that is collected is related to incident prevention, as opposed to collecting information on the effects of an incident once it occurs (H. van Poelje: personal communication, November 30, 1995).

Project management for EGIG is based in the Netherlands (at N.V. Nederlandse Gasunie), and so it is appropriate to further discuss the Dutch regulations. As explained in Chapter 1, the Dutch have established four area classes, with effect and building distances determined through a risk analysis. One of the possible ways in which risk
contours can be determined involves establishing the pipeline failure frequencies, determining the probability of ignition and establishing the dose-effect relationship (H. van Poelje: personal communication, September 30, 1996). It should be noted that there is no internationally agreed method for risk calculation. While the Dutch definition of individual risk is based on fatalities, this may not be the basis used by other countries. Another basis which other countries use is casualties.

### 2.5.3 Research Performed by the British Health and Safety Executive

In their research report to the HSE, Hill and Catmur (31) compare individual and societal risks for different high pressure pipelines including methane, oxygen, spiked crude oil, NGL (natural gas liquids), gasoline, ethylene and ammonia. The authors define two criteria for evaluating risk. "Fatality" is defined as an impact likely to cause death immediately or after an extended duration, while "dangerous dose or worse" is defined to include:

- $\quad$ Severe distress to many persons;
- A substantial proportion of persons affected require medical attention;
- Some persons are seriously injured, requiring prolonged treatment; and - Any highly susceptible people might receive a fatal injury.

Individual risk transects and societal $\mathrm{F} / \mathrm{N}$ curves were prepared. As defined in the report, The individual risk transects present the probability (at varying distances) that an individual could incur a fatal injury or dangerous dose due to an accidental release from a pipeline. $\mathrm{F} / \mathrm{N}$ curves present societal risk as the probability of N or more fataiities due to a pipeline release.

The authors determine that all pipelines except oxygen and small diameter low pressure methane pipelines had individual risk levels greater than one in a million per year. Additionally, the greatest number of multiple fatalities were associated with methane, NGL and ethylene pipelines.

In a second research report prepared for the HSE, Hockey and Rew (32) review the factors that influence the response of humans to thermal radiation. The report also discusses and compares various methodologies for determining the effects of thermal radiation from process fires. The authors indicate that most risk assessment methodologies consider the speed of escape, delay before escape begins and the distance traveled to reach shelter as factors influencing the probability of fatality.

### 2.5.4 Research Performed in Canada

Based on information provided by the TSB of Canada (personal communication, February 28, 1997), a private company in Canada is developing, in coordination with the Major Industrial Co-ordinating Council of Canada (MIACC), a document intended for use by government officials during the planning of pipeline routes and development. The document defines a Response Planning Area (RPA) to be an area where public safety
issues (such as emergency response and preparedness) should be addressed. The RPA is developed using consequence modeling, and is based on the following major factors:

- Source parameters (e.g., pipeline diameter and flow, and the properties of the transported material);
- Transport parameters which are used to evaluate the migration of material or energy from the pipeline; and
- Damage criteria

The RPAs are intended to be presented as graphs (for example, distance to a specified damage criteria as a function of pipeline diameter and pressure).

In addition to the aforementioned information, the TSB of Canada likewise provided (personal communication, February 28, 1997) information on guidelines, developed by the Canadian Standards Association, for performing risk analyses on pipelines in Canada. As stated in these guidelines, the intent is to:
"(a) identify the role of risk analysis within the context of an overall risk management process;
(b) set out standard terminology that is consistent with existing Canadian standards in the field of risk management;
(c) identify in general terms the components of the risk analysis process, the associated data requirements, and the requirements for documentation and records; and
(d) where applicable, provide reference to methodological guidelines for risk analysis".

### 2.5.5 Research Performed by NJIT

The NJIT Institute for Transportation (33) has reviewed NTSB Pipeline Accident Reports (PARs) for the purpose of determining the burn radius associated with natural gas pipeline failures involving fires. The authors acknowledge that there are many factors that contribute to the development of a burn radius (e.g., rupture geometry, wind conditions, terrain geometry, duration of release and quantity of material released, etc.). However, the authors maintain that it is still possible to determine an upper bound on the extent of the burn radius. The upper bound was found to be a 92 foot radius at an incident operating pressure of 260 psig, and a 610 foot radius at a pressure of 987 psig.

The authors also found a close correlation between the occurrence of a natural gas pipeline explosion and fire, and its resultant burn radii. Fire damage is a function of the radiant energy produced by a fire, and the authors indicate that after an explosion there is less natural gas available to produce radiant energy. Therefore, the authors conclude that the burn radius associated with explosions and fires is less than that associated with only fires for similar natural gas pipeline operating conditions.

### 2.6 Conclusions from Literature Review

The prior research reviewed and cited herein is predominantly concerned with predicting which sections of a pipeline are most vulnerable to damage and predicting the loss of product during an actual pipeline rupture (i.e., blowdown), rather than directed to concerns associated with establishment of adequate safe separation distances for the public.

The above research emphasis probably exists because the risk of injury and fatality to the general public related to natural gas pipeline ruptures has been documented by reviewing agencies to be minimal compared to other incidents which human beings may encounter (i.e., due to falls, drowning, poisoning, fires and burns, etc.) during their normal lifetime.

The modeling of product loss during a pipeline rupture, which is an important factor related to the establishment of a safe separation distance therefrom, is, as indicated in a number of research papers reviewed, difficult to simulate. This is due to disparity amongst researchers as to whether blowdown in the pipeline should be modeled as an adiabatic or isothermal condition; whether the fluid is viscous or nonviscous; whether the flow regime is isentropic or nonisentropic and consequently whether the natural gas should be treated as a perfect gas or a real fluid. Further, the prior research suggests that the gas may change from one physical state to another within the pipeline during the course of the same blowdown event.

In addition to the above-noted complexities in attempting to accurately simulate the fluid dynamics within a pipeline during a blowdown situation, there are a number of other factors which directly influence the required safe separation distance that should be established between a specific natural gas pipeline and its neighboring potentially impacted public at large. These factors include: the pipeline diameter; operating pressure; location of and time needed for emergency closure of the closest valves located upstream and downstream of the rupture; the interconnections (if any) with other pipelines in the vicinity of the rupture; the physical orientation (i.e., the nature and
location of the break) of the rupture; the depth of cover over the pipeline; and the atmospheric conditions (e.g. wind velocity, wind direction, etc.) in the immediate vicinity of the pipeline at the time of rupture. In attempting to predict, with the use of a generalized mathematical model, the impacts of all these variables both independently and jointly, it appears that an accurate single point estimate of associated safe separation distance would be exceedingly difficult if not impossible to develop.

The available hazard assessment models reviewed herein have been developed primarily for utilization by private industry and the Federal government, and are designed for generalized hazard assessment purposes. The models can be extremely expensive to purchase, and are, in many instances, proprietary in nature. Furthermore, specific documentation associated with the nature of the model development, such as the structure of the computer codes, is lacking in the published material associated with the above-noted models. As such, attempts to extract meaningful information related to this thesis have resulted in limited success.

The literature review also revealed that the available database of information associated with actual natural gas pipeline accident occurrences in the United States is limited. For example, it was previously indicated that the NTSB investigates and reports upon pipeline accidents in the United States. However, the NTSB does not investigate all pipeline accidents; the criteria that triggers an actual NTSB investigation relates to the extent of property damage and whether injuries or fatalities result from the accident. The results of the investigations that are performed are published as PARs. Also, a review of the available PARs by this author indicates that there is an inconsistency in the type of
information gathered by the NTSB for each accident. For example, the PARs do not consistently report the total volume of gas lost in an accident, nor the location of the closest valves upstream and downstream of the pipeline rupture. Furthermore, supporting information for the more dated PARs (in the PAR dockets in Washington, DC) is periodically destroyed.

The Commodity Pipeline Occurrence Reports prepared by the Canadian TSB are similar to the PARs in that there are also inconsistencies in the extent of information contained in these reports. Further, the Canadian TSB investigates an extremely small number of accidents (i.e., approximately one percent of all pipeline accidents in Canada).

Because of all the aforementioned reasons, it is concluded herein that an approach to advance the state-of-the-art in the discipline of pipeline risk analysis is to develop a reliable estimation technique to conservatively predict safe separation distances to be articulated between the public and the ruptured natural gas pipeline. To aid in this development, the literature review indicates that extensive work has been performed in determining turbulent jet flame characteristics. In addition, the previously noted point source model can be utilized to calculate the heat flux generated at various distances from industrial flares. As will be seen in the following chapters, this model can be applied to the evaluation of safe separation distances which can be considered by regulatory agencies to protect the quality of life of residents located in proximity to natural gas pipelines.

## CHAPTER 3

## OBJECTIVE OF THE DISSERTATION AND

 INTRODUCTION TO THE DISSERTATION RATIONALEAs noted in Chapter 2 in the literature search and the analysis thereof, there is considerable complexity involved in attempting to mathematically model the dynamics associated with a natural gas pipeline rupture. This complexity arises due to the numerous variables involved in the process as well as the varying assumptions made by previous investigators as to the nature of the physical state of the natural gas during a rupture condition. In consideration of the above, proposed herein is a simplified approach for estimating safe separation distances based upon assumptions that the damage from a pipeline rupture is primarily due to the thermal radiation produced by the ignited gas behaving like a vertical jet flame, and that the major variables associated with the burn radii (i.e., impact area) resulting from a pipeline rupture are the size of the pipeline and its operating pressure (which directly affects the mass flow rate in the pipeline). The appropriateness of this approach in providing results of an accuracy suitable for regulatory agencies to utilize for establishing zoning guidelines is confirmed in this thesis by comparing results from the aforementioned model with actual burn radii found in a limited number of accident investigations conducted by the NTSB in which sufficient data was obtained to verify the subject model herein.

As indicated above, an assumption is made that the damage from a pipeline rupture is primarily due to the thermal radiation produced by the ignited gas. Therefore, a safe separation distance is defined as the distance beyond which a pre-established level of thermal radiation damage will not likely occur. Other damage (such as projectile damage from pipeline fragments or damage due to overpressures from explosions) is not considered. This assumption is consistent with the results from investigations of actual pipeline ruptures, in which damage was found to be caused primarily from the fire.

It was also previously mentioned that the escaping gas is assumed to behave, once ignited, like a vertical jet flame. A release from a pressurized system like a pipeline can produce other scenarios such as dispersion of the unignited gas, formation of a fireball, development of a flash fire or a vapor cloud explosion (20, 31, 34). Furthermore, the rupture orientation may be such that a flame jet, if it exists, may not be truly vertical.

Assuming that all of the above scenarios can occur increases the number of variables to be considered, in that the probabilities of each scenario happening (either alone or in combination with other scenarios) must then be determined. Furthermore, should these other scenarios occur, there is no certainty that they will contribute significantly to the overall thermal radiation damage. For example, the dispersion of unignited gas would not produce thermal radiation damage. Vapor cloud explosions can produce damage through the generation of pressure waves. However, according to the Center for Chemical Process Safety (CCPS), the following conditions need to be present for a vapor cloud explosion to occur (34):

- The released material must be flammable and at an appropriate temperature and pressure.

A suitably-sized cloud of the escaping gas must be present before ignition. CCPS reports that ignition delays of 1 to 5 minutes are the most probable for generating clouds which can explode.

- A sufficient portion of the cloud must be within the flammable range of the material.
- The rate of flame propagation within the cloud determines the severity of blast effects. High flame speeds can produce high overpressures.

As indicated previously, thermal radiation is the primary cause of damage in natural gas pipeline ruptures. Since vapor cloud explosions are produced only under a set of very specific circumstances and since blast effects are not the primary cause of damage, this scenario is not expected to routinely occur.

With regard to the development of a flash fire, CCPS (34) indicates that very little information is currently available concerning the thermal radiation produced. CCPS also states that thermal radiation hazards from burning vapor clouds are considered less significant than blast effects, and that combustion associated with a flash fire lasts no more than a few tens of seconds. In addition, Hockey and Rew (32) indicate that for the
calculation of hazard ranges, it is assumed that people in contact with the flame are fatally burned while those at a distance would not be seriously affected. It should be noted that the fire dimensions will depend on the shape of the vapor cloud, which itself is dependent upon many variables. Given these limitations, the modeling of flash fires for the determination of safe separation distances is not considered in the dissertation.

Although fireballs are typically associated with BLEVEs (Boiling Liquid Expanding Vapor Explosions), they can also happen as a result of a pipeline rupture. CCPS (34) lists the properties of firebails which have the greatest influence on thermal radiation. These include:

- Fireball diameter as a function of time and maximum diameter;
- Height of fireball center above its ignition position as a function of time elapsed after liftoff;
- Fireball surface emissive power; and
- Total combustion duration.

Although fireballs produce the highest radiation intensity, these events can be assumed to last only 10-30 seconds (32). Formulas for fireball diameter, duration and hazard distances have been published (34) which are functions of the mass of the fuel. However, in the case of a pipeline rupture the mass of fuel involved in a fireball is difficult to predict since the release rate varies with time. Due to the limitations and numerous variables involved, fireball modeling is not attempted in the dissertation.

In order to model a vertical flame jet emanating from a pipeline rupture, the point source model will be used. As indicated previously, the point source model was used by Oenbring and Sifferman (24) to calculate the heat flux from a flare. Since the damage incurred by an object receiving radiant heat from a flame is dependent upon the heat flux, the point source model can be used to relate thermal radiation damage to the distance from the ruptured pipeline. The model is dependent upon the total heat content of the escaping gas (which is a function of the gas release rate) and the height of the flame (for location of the point source). Therefore, the following chapter of the dissertation will develop relationships for flow rate and flame height which will be used in conjunction with the point source method for predicting safe separation distances.

## CHAPTER 4

## DEVELOPMENT OF METHOD TO DETERMINE SAFE SEPARATION DISTANCES

This chapter presents the methodology which is used to estimate safe separation distances based on pipeline diameter and incident operating pressure. The equations for determining heat flux, gas flow rate and flame height are presented and combined to produce a single safe separation distance equation. This equation is then used in the preparation of charts for predicting safe separation distances.

### 4.1 Calculation of Burn Radii

As indicated in Chapter 3, the majority of damage resulting from a pipeline rupture is caused by thermal radiation. Therefore, the safe separation distance from a pipeline can be defined in terms of the distance needed to protect against a specified heat flux. This specified heat flux will produce an area of thermal radiation damage in the vicinity of the pipeline which can be estimated by calculating the burn radius. In Chapter 3, the assumption was made that a pipeline rupture produces a vertical jet flame. An assumption is now made that this vertical flame will radiate equivalently in all directions so that the burn area will be symmetric about the rupture. The burn area can then be described by a circle, with the radius of the circle being the burn radius. The following section describes how the burn radius can be determined using the point source method.

### 4.1.1 Use of the Point Source Method

As explained in Chapter 3, the point source method forms the basis for the estimation of safe separation distances. In Chapter 2, the following equation was presented from the work of Oenbring and Siffferman (24):

$$
\begin{equation*}
\mathrm{K}=\mathrm{FQ} /\left(4 \pi \mathrm{D}^{2}\right) \tag{2-33}
\end{equation*}
$$

where
$\mathrm{K}=$ Radiation heat flux from a flame (Btu/hr $\mathrm{ft}^{2}$ )
$F=$ Fraction of total heat radiated
$\mathrm{Q}=$ Total heat content of the flared gas (Btu/hr)
$\mathrm{D}=$ Distance from point source to receptor (ft)

The authors provide the source for Equation 2-33 as being the American Petroleum Institute (API) document API RP-521. In a later version of this document (35), API provides the revised equation:

$$
\begin{equation*}
\mathrm{D}=(\tau \mathrm{FQ} /(4 \pi \mathrm{~K}))^{1 / 2} \tag{4-1}
\end{equation*}
$$

where
$\tau=$ the fraction of thermal radiation transmitted through the atmosphere
Application of the point source method is shown in Figure 4.1, where ignition of escaping gas from a pipeline rupture results in a flame of height " $H$ ". The point source is placed in the center of the flame at $H / 2$, and the burn radius $(B R)$ is found from the Pythagorean Theorem:

$$
\begin{equation*}
B R=\left\{D^{2}-(H / 2)^{2}\right\}^{1 / 2} \tag{4-2}
\end{equation*}
$$

By inserting Equation 4-1 into Equation 4-2, the following relationship is obtained:

$$
\begin{equation*}
\mathrm{BR}=\left\{(\tau \mathrm{FQ} /(4 \pi \mathrm{~K}))-(\mathrm{H} / 2)^{2}\right\}^{1 / 2} \tag{4-3}
\end{equation*}
$$

This is the basic form of the burn radius equation. The burn radius is function of the transmissivity of the thermal radiation through the atmosphere, the fraction of total heat radiated, the heat content of the escaping gas and the specified heat flux (or level of damage).

Based on information provided by various researchers for methane (20), the value of F can be reasonably estimated to be 0.2. Furthermore, APl (35) discusses the work of Brzustowsi and Sommer (36) for calculating transmissivity of thermal radiation through the atmosphere. The authors indicate that $\tau$ can be found from the following equation:

$$
\begin{equation*}
\tau=0.79(100 / \mathrm{r})^{1 / 16}(100 / \mathrm{D})^{1 / 16} \tag{4-4}
\end{equation*}
$$

where
$r=$ relative humidity, $\%$
$\mathrm{D}=$ distance to flame, feet

The authors caution that Equation 4-4 is "strictly applicable" when the flame is radiating at a temperature of $2240^{\circ} \mathrm{F}$, relative humidity is greater than $10 \%$, dry bulb temperature is $80^{\circ} \mathrm{F}$ and the distance from the flame is greater than 100 feet but less than 500 feet. However, the authors indicate that an "order of magnitude" estimate can be made under a wider range of conditions. Since the results from pipeline accident investigations typically show damage extending several hundred feet from the rupture, an estimate for $\tau$ is found by assuming a relative humidity of $50 \%$ and a distance to the flame of 500 feet (both of which are reasonable values):

$$
\tau=0.79(100 / 50)^{1 / 16}(100 / 500)^{1 / 16}=0.746
$$

Inserting the values of $\tau$ and $F$ into Equation 4-3, the equation for the burn radius becomes:

$$
\begin{align*}
\mathrm{BR} & =\left\{(0.746(0.2) \mathrm{Q} /(4 \pi \mathrm{~K}))-(\mathrm{H} / 2)^{2}\right\}^{1 / 2} \\
& =\left\{(0.011873 \mathrm{Q} / \mathrm{K})-(\mathrm{H} / 2)^{2}\right\}^{1 / 2} \tag{4-5}
\end{align*}
$$

The total heat content of the escaping gas $(Q)$ in $B t u / h r$ can be found by multiplying the heat content of natural gas ( $1,000 \mathrm{Btu} / \mathrm{scf}$ ) by the volumetric flow rate of the escaping gas (scf/hr):

$$
\begin{equation*}
\mathrm{Q}=1,000\left(\mathrm{~V}^{\prime}\right) \tag{4-6}
\end{equation*}
$$

where
$\mathrm{V}^{\prime}=$ Volumetric flow rate ( $\mathrm{scf} / \mathrm{hr}$ )

If Equation 4-6 is combined with Equation 4-5, the expression for the burn radius becomes:

$$
\begin{equation*}
\mathrm{BR}=\left\{\left(11.873\left(\mathrm{~V}^{\prime}\right) / \mathrm{K}\right)-(\mathrm{H} / 2)^{2}\right\}^{1 / 2} \tag{4-7}
\end{equation*}
$$

The following sections of this chapter will provide discussions of the three variables in equation 4-7: heat flux $(\mathrm{K})$; gas flow rate $\left(\mathrm{V}^{\prime}\right)$; and flame height $(\mathrm{H})$.

### 4.1.2 Heat Flux Values

Examples of heat flux values corresponding to specific consequences are provided in Table 4.1. These values were obtained from a review of the literature. A comprehensive listing of heat flux values can be found in Appendix A.

## Table 4.1

Heat Flux Values

| Heat Flux |
| :--- |
| Btu/hr ft |

$222^{2} *$
** Calculated using the relationship $1 \mathrm{Btu} / \mathrm{hr}_{\mathrm{ft}^{2}}=3.1546 \mathrm{~W}$ atts/square meter $\left(\mathrm{W} / \mathrm{m}^{2}\right)$.

From several of the above examples, it can be seen that the level of thermal radiation damage may not only depend on the intensity of the heat flux, but also on the length of time that the receptor is receiving that heat flux. For example, at a heat flux
intensity of $9,985 \mathrm{Btu} / \mathrm{hr} \mathrm{ft}^{2}$, spontaneous ignition of wooden buildings occurs after a few minutes. At $555 \mathrm{Btu} / \mathrm{hr}_{\mathrm{ft}^{2}}$, the pain threshold is reached after 60 seconds.

Therefore, a safe separation distance will be considered to afford protection from a certain level of heat flux for a specific time period. If that time period is exceeded, damage may occur.

In order to estimate a safe separation distance, a level of protection is chosen, such as protecting wooden buildings from spontaneous ignition for a few minutes. The corresponding heat flux is found and inserted into Equation 4-7 as the appropriate K -value. The procedure described in the next section is used to determine an appropriate gas flow rate.

### 4.1.3 Determination of Gas Flow Rate

In order to determine the flow rate of gas escaping from a pipeline rupture, a variation of Equation 2-15 will be used. The use of Equation 2-15 for determining the gas rate from full bore ruptures is similar to the procedure used by Wilson (11), in which Equation 2-5 (for choked conditions) was used to determine the initial mass release rate from ruptures. While choked conditions would likewise be expected with high pressure pipeline punctures and blowdowns, the use of Equation 2-15 provides a convenient means of estimating gas flow rate knowing only the pipeline diameter and an upstream pressure. Recall that Equation 2-15, as presented below, is applicable for use during pipeline punctures and blowdowns:

$$
\begin{equation*}
\mathrm{Q}=\mathrm{D}^{2} \mathrm{P}_{1} \tag{2-15}
\end{equation*}
$$

As described previously, " $Q$ " represents the volume of gas in Mcf/hr at a pressure of 14.9 psi and $60^{\circ} \mathrm{F}$ with a specific gravity of 0.60 . In other words, Q provides the flow rate at approximately standard conditions. To use Equation 2-15, "D" (the diameter of the nipple or orifice in inches) is considered to be the nominal diameter of the pipeline for a worst-case blowdown scenario. Since " $\mathrm{P}_{1}$ " is the absolute pressure (in psi ) at a point upstream from the opening, this variable will represent the pre-rupture incident operating pressure at the location of the rupture.

By examining the rate of gas lost through actual pipeline ruptures and comparing this rate with values obtained using Equation 2-15, an evaluation can be made as to whether a modifying factor must be applied to Equation 2-15. Table 4.2 presents eight pipeline accidents investigated and reported upon by the NTSB and the Canadian TSB in which the following parameters are known: pipeline diameter, incident operating pressure, rupture isolation time, and volume of gas lost.

## Table 4.2

Pipeline Rupture Parameters

| Accident <br> Report <br> Number | Reference <br> Number | Pipeline <br> Diameter <br> Inches | Incident <br> Operating <br> Pressure <br> psia | Isolation <br> time <br> hours | Total Volume <br> of Gas <br> Lost, scf |
| :--- | :--- | :--- | :--- | :--- | :--- |
| P90H0606 | 41 | 12.75 | 696.7 | 2.75 | $3.78 \times 10^{7}$ |
| 83-02 | 42 | 20 | 834.7 | 1.42 | $4.68 \times 10^{7}$ |
| P91H0041 | 43 | 20 | 933.7 | 0.75 | $3.13 \times 10^{6}$ |
| 79-FP006 | 44 | 30 | 574.7 | 2.83 | $2.01 \times 10^{8}$ |
| P90H1006 | 45 | 30 | 726.7 | 0.58 | $8.73 \times 10^{7}$ |
| 95-01 | 3 | 36 | 984.7 | 2.50 | $2.97 \times 10^{8}$ |
| P94H0036 | 46 | 36 | 1014.7 | 0.63 | $1.48 \times 10^{8}$ |
| P94H0003 | 47 | 42 | 1221.7 | 6.67 | $3.52 \times 10^{8}$ |

As described in Chapter 2, critical conditions produce the maximum gas flow rate.
Therefore, it will be assumed that, using the incident operating pressure, Equation 2-15 (or Equation 2-15 with a modifying factor) will predict a maximum initial flow rate.

Furthermore, this maximum flow rate will decrease as the pipeline is depressurized. The maximum flow rate will be identified as $\mathrm{V}_{\mathrm{CALC}}$.

Another flow parameter that can be calculated is the average flow rate, $V_{A \vee G}$. The average flow rate represents the total volume of gas that is lost divided by the time required to isolate the ruptured pipeline segment. Since $V_{A V G}$ represents an average, or constant value of flow, the flow rate remains the same from initiation of rupture to isolation of the pipeline.

Table 4.3 presents the results of computations of $V_{A V G}, V_{\text {CALC, }}$, and the ratio $\mathrm{V}_{\mathrm{AVG}} / \mathrm{V}_{\mathrm{CALC}}$.

## Table 4.3

Flow Rate Parameters

| Accident |  |  |  |
| :---: | :---: | :---: | :---: |
| Report | $\mathrm{V}_{\mathrm{AVG}}$ | $\mathrm{V}_{\mathrm{CALC}}$ | $\mathrm{V}_{\mathrm{AVG}} / \mathrm{V}_{\mathrm{CALC}}$ |
|  |  | scf/hr | Ratio |
| P90H0606 | $1.37 \times 10^{7}$ | $1.13 \times 10^{8}$ | 0.12 |
| 83-02 | $3.30 \times 10^{7}$ | $3.34 \times 10^{8}$ | 0.10 |
| P91H0041 | $4.17 \times 10^{6}$ | $3.73 \times 10^{8}$ | 0.01 |
| 79-FP006 | $7.10 \times 10^{7}$ | $5.17 \times 10^{8}$ | 0.14 |
| P90H1006 | $1.51 \times 10^{8}$ | $6.54 \times 10^{8}$ | 0.23 |
| 95-01 | $1.19 \times 10^{8}$ | $1.28 \times 10^{9}$ | 0.09 |
| P94H0036 | $2.35 \times 10^{8}$ | $1.32 \times 10^{9}$ | 0.18 |
| P94H0003 | $5.28 \times 10^{7}$ | $2.16 \times 10^{9}$ | 0.02 |

For the accidents listed above, $\mathrm{V}_{\mathrm{AVG}}$ is found to vary between $1 \%$ and $23 \%$ of $V_{\text {CALC }}$, with an average of $11 \%$. Since during an actual pipeline rupture the initial mass flow rate will be very high, the flow rate would not be expected to have a value lower than $V_{A V G}$. Furthermore, if the assumption is made that the actual flow rate will not be greater than $\mathrm{V}_{\mathrm{CALC}}$, then the actual initial flow rate would be expected to lie between $\mathrm{V}_{\mathrm{AVG}}$ and $\mathrm{V}_{\mathrm{CALC}}$. Therefore, assume that the initial flow rate is found at the midpoint between $0.11 \mathrm{~V}_{\text {CALC }}$ (i.e., $\mathrm{V}_{A V G}$ ) and $\mathrm{V}_{\text {CALC, }}$ which is equal to $0.56 \mathrm{~V}_{\text {CALC. }}$ This result is very similar to the result found by Wilson and expressed as Equation 2-11, where the initial mass release rate was found to be $50 \%$ of the value found using Equation 2-5 (12).

Inserting the maximum initial flow rate into the point source equation would not accurately reflect thermal radiation conditions, since the heat fiux at a given receptor location will decrease with the decreasing gas flow. Therefore, assume that a representative gas flow rate to use in the point source equation exists between 0.56 $\mathrm{V}_{\mathrm{CALC}}$ and $\mathrm{V}_{\mathrm{AVG}}$ (or $0.11 \mathrm{~V}_{\mathrm{CALC}}$ ). The midpoint of this range is $0.34 \mathrm{~V}_{\mathrm{CALC}}$, which then leads to the following equation for determining flow rate:

$$
\begin{equation*}
V^{\prime}=(1,000)(0.34) D^{2} p \tag{4-8}
\end{equation*}
$$

where
$\mathrm{V}^{\prime}=$ Gas flow rate, scf/hr
$D=$ Pipeline Diameter, inches
$P=$ Incident Operating Pressure, psia
It should be noted that multiplication by 1,000 in Equation $4-8$ converts $\mathrm{Mcf} / \mathrm{hr}$ to $\mathrm{scf} / \mathrm{hr}$.

By inserting Equation 4-8 into Equation 4-7, the following relationship is
obtained:

$$
\begin{align*}
\mathrm{BR}= & \left\{\left(11.873(1,000)(0.34) \mathrm{D}^{2} \mathrm{P} / \mathrm{K}\right)-(\mathrm{H} / 2)^{2}\right\}^{1 / 2} \\
& \left.=\left\{(4,036.82) \mathrm{D}^{2} \mathrm{P} / \mathrm{K}\right)-(\mathrm{H} / 2)^{2}\right\}^{1 / 2} \tag{4-9}
\end{align*}
$$

The burn radius is now expressed as a function of the pipeline diameter, operating pressure, heat flux and the flame height. In the next section of this chapter, an expression for the flame height will be developed which can be inserted into Equation 4-9.

### 4.1.4 Determination of Flame Height

In a manner similar to the method for determining the expression for gas flow rate, the information obtained from actual pipeline accidents can be used to estimate flame height.

Recalling the Hawthorne Equation that was provided in Chapter 2:

$$
\begin{equation*}
\mathrm{L} / \mathrm{D}=\left(5.3 / \mathrm{C}_{\mathrm{T}}\right)\left\{\left(\mathrm{T}_{\mathrm{F}} /\left(\alpha_{\mathrm{T}} \mathrm{~T}_{\mathrm{N}}\right)\right)\left(\mathrm{C}_{\mathrm{T}}+\left(1-\mathrm{C}_{\mathrm{T}}\right)\left(\mathrm{M}_{\mathrm{S}} / \mathrm{M}_{\mathrm{N}}\right)\right\}^{1 / 2}\right. \tag{2-16}
\end{equation*}
$$

where, in dimensions of mass $(M)$, length $(\mathrm{L})$, time $(\mathrm{t})$ and temperature $(\mathrm{T})$, the various terms are:
$L=$ visible flame length ( $L$ )
$D=$ nozzle diameter (L)
$T_{\mathrm{F}}=$ adiabatic flame temperature (absolute T )
$\mathrm{T}_{\mathrm{N}}=$ temperature of nozzle fluid (absolute T )
$\mathrm{M}_{\mathrm{S}}=$ molecular weight of surrounding fluid (unitless)
$\mathrm{M}_{\mathrm{N}}=$ molecular weight of nozzle fluid (unitless)
$\mathrm{C}_{\mathrm{T}}=$ mol fraction of nozzle fluid in the unreacted stoichiometric mixture (decimal)
$\alpha_{T}=$ mols of reactants/mols products for the stoichiometric mixture (decimal)

Equation 2-16 is applicable to vertical flame jets in still air. Since the assumption was made that the flame jet from a pipeline rupture will be vertical, the equation can be used in the evaluation of flame height.

Using the values of $\mathrm{C}_{\mathrm{T}}=0.091, \alpha_{\mathrm{T}}=1$ and $\mathrm{T}_{\mathrm{F}} / \mathrm{T}_{\mathrm{N}}=7.4$ for methane (20) in Equation 2-16, along with a molecular weight of 29.0 (48) for air and a molecular weight of 16.042 for methane (8), the following expression is obtained:

$$
\begin{gather*}
\mathrm{L} / \mathrm{D}=(5.3 / 0.091)\{(7.4 / 1)(0.091+(1-0.091)(29.0 / 16.042))\} 1 / 2 \\
=208.6 \tag{4-10}
\end{gather*}
$$

Therefore, for a given gas, the flame length is directly proportional to the jet (or nozzle) diameter. This observation can be applied to the NTSB pipeline accident data, where estimation of flame heights are provided. If the assumption is made that the jet diameter is equal to the pipeline diameter for a full bore rupture, the following is obtained:

Table 4.4
Flame Height Data
$\left.\begin{array}{lllllll}\begin{array}{lll}\text { Accident } \\ \text { Report }\end{array} & \begin{array}{lll}\text { Reference } \\ \text { Number }\end{array} & \begin{array}{l}\text { Diameter } \\ \text { (D) }\end{array} & \begin{array}{l}\text { Diameter } \\ \text { (D/12) }\end{array} & \begin{array}{l}\text { Reported Flame } \\ \text { Height }\end{array} & & \begin{array}{l}\text { Feet } \\ \text { (H). Feet }\end{array}\end{array}\right)$

The average $H /(D / 12)$ is 147 , so assume that this ratio can be used to estimate flame height. Therefore, the following relationship is obtained:

$$
\begin{align*}
\mathrm{H} & =147(\mathrm{D} / 12) \\
& =12.25(\mathrm{D}) \tag{4-11}
\end{align*}
$$

where
$\mathrm{D}=$ Pipeline diameter (in)
While it is understood that the calculated L/D value (208.6) is of a similar order of magnitude to the average $\mathrm{H} / \mathrm{D}$ value (147), it is this average $\mathrm{H} / \mathrm{D}$ value that will be used in subsequent calculations. This is justified since the original work of Hawthorne (21) used small diameter nozzles to produce the jet instead of full pipeline diameters.

By inserting equation 4-11 into Equation 4-9, the final burn radius equation is found:

$$
\begin{align*}
\mathrm{BR}= & \left.\left\{(4,036.82) \mathrm{D}^{2} \mathrm{P} / \mathrm{K}\right)-(12.25 \mathrm{D} / 2)^{2}\right\}^{1 / 2} \\
& =\mathrm{D}\{(4,036.82) \mathrm{P} / \mathrm{K})-37.52\}^{1 / 2} \tag{4-12}
\end{align*}
$$

where
$\mathrm{BR}=$ Burn Radius $(\mathrm{ft})$
$\mathrm{D}=\operatorname{Pipeline}$ Diameter (in)
P = Incident operating pressure (psia)
$\mathrm{K}=$ Heat flux (Btu/hr $\mathrm{ft}^{2}$ )

Equation 4-12 provides a means by which the burn radius (and hence the safe separation distance) can be found knowing only the pipeline diameter, incident operating pressure and the level of damage (i.e., heat flux) to be considered. Since pipeline operating pressures are typically specified as gauge pressures, Equation 4-12 can be modified for application to gauge pressures by substituting the quantity ( $\mathrm{P}^{\prime}+14.7$ ) for P , where $P^{\prime}$ is the incident operating pressure in psig.

Using Equation 4-12, charts can be developed for estimating the burn radius. The method for constructing these charts is presented in the next section.

### 4.1.5 Construction of Charts to Predict Safe Separation Distances

The following procedure is used to construct charts for the rapid estimation of safe separation distances. The first step involves deciding what degree of thermal radiation damage to consider. For example, the damage might be spontaneous ignition of wooden buildings after a few minutes exposure to the ignited gas. Protection for a few minutes may allow enough time for emergency responders to arrive at the scene and initiate protective measures (such as watering down the buildings). Based on the information in both Appendix A and Table 4, the heat flux corresponding to the specified level of damage is $9,985 \mathrm{Btu} / \mathrm{hr} \mathrm{ft}^{2}$. This value is inserted into Equation $4-12$ for K :

$$
\begin{equation*}
\mathrm{BR}=\mathrm{D}\{(4,036.82) \mathrm{P} / 9,985)-37.52\}^{1 / 2} \tag{4-13}
\end{equation*}
$$

Equation 4-13 is an expression of the burn radius as a function of only diameter and incident operating pressure. For various pipeline diameters, graphs are then constructed of the burn radius (on the $y$-axis) and the incident operating pressure (on the $x$-axis). In


Figure 4.2 - Example of a Burn Radius Chart for a 36" Diameter Pipeline, $K=9,985 \mathrm{Btu} / \mathrm{hr} \mathrm{sq}$. ft .
the example, a pipeline diameter of $36^{\prime \prime}$ can be used with incident operating pressures in the range of 575 psia to 1,200 psia to construct a chart similar to the one shown in Figure 4.2. Once the graph is completed, it can be used either to determine a safe separation distance given a specified incident operating pressure, or to determine the incident operating pressure required to maintain a specified safe separation distance.

Charts for estimating safe separation distances (or burn radii) are provided in Appendix B for heat fluxes (from Appendix A) of $3,962 \mathrm{Btu} / \mathrm{hr} \mathrm{ft}^{2}$ (piloted ignition of wood); $6,340 \mathrm{Btu} / \mathrm{hr} \mathrm{ft}^{2}$ (blistering of bare skin in 4 seconds and $1 \%$ lethality in 20 seconds); $9,510 \mathrm{Btu} / \mathrm{hr} \mathrm{ft}^{2}$ (causes third degree burns in 30 seconds); and $9,985 \mathrm{Btu} / \mathrm{hr} \mathrm{ft}^{2}$ (spontaneous ignition of wooden structures after a few minutes). The charts have been developed for pipeline diameters of $14^{\prime \prime}, 16^{\prime \prime}, 18^{\prime \prime}, 20^{\prime \prime}, 24^{\prime \prime}, 30^{\prime \prime}$ and $36^{\prime \prime}$, with incident operating pressures in the range of 575 psia to $1,200 \mathrm{psia}$. It should be noted that the charts do not consider that portion of the heat flux due to solar radiation. An accurate value of the solar heat flux would be dependent on factors such as the weather conditions, the time of day and the time of year. Since the solar heat flux amounts to only a few hundred Btu/hr $\mathrm{ft}^{2}$ while the non-solar heat flux is several thousand $\mathrm{Btu} / \mathrm{hr}_{\mathrm{ft}}{ }^{2}$, omission of this factor will not significantly affect the results.

Equation 4-12 will be used in the next chapter for comparing (for verification purposes) predicted values of burn radii using the model developed herein to information from actual pipeline accidents and to the results of prior research.

## CHAPTER 5

## COMPARISON OF METHOD TO PIPELINE ACCIDENT DATA AND TO PREVIOUS RESEARCH

The previous chapter presented Equation 4-12 as a means by which a burn radius can be determined through knowledge of a pipelines's diameter, incident operating pressure and the level of thermal radiation damage. In this chapter, the utilization of Equation 4-12 will be evaluated by first comparing the results it produces to data from two highly documented pipeline accidents that occurred in the United States. These accidents (located in Edison, New Jersey and Lancaster, Kentucky) were selected for this detailed analysis since these were the only two accidents for which extensive descriptions of the thermal radiation damage could be obtained from the NTSB. Additionally, an attempt is made to compare the results from Equation 4-12 with other pipeline accidents in which less extensive information is reported.

The second part of this chapter will compare the results of Equation 4-12 to separation distance calculations that have been developed by other countries (i.e., Great Britain and the Netherlands) through the use of the principles of hazard analysis.

### 5.1 Comparison to Pipeline Accident Data

This section provides a comparison of burn radii obtained from actual pipeline rupture events to corresponding estimates developed through use of Equation 4-12. Since review of the pipeline accident information reveals that burn patterns are not truly symmetric (due to numerous factors including variation in wind velocity and direction during a rupture condition), the actual burn radii will be defined conservatively as the farthest horizontal distance from the rupture at which the specified level of thermal radiation damage is observed. Since the point source method assumes a circular (i.e. symmetric) burn pattern around the rupture, there may be locations around the rupture for which Equation 4-12 overestimates the extent of thermal radiation damage based upon the above noted assumption. However, comparison of the results obtained from Equation 4-12 to the maximum horizontal burn distances is an approach which provides a level of confidence that the estimated burn radii should approximate, and not be significantly less than, the actual burn radii (which would be undesirable in actual practice).

The following section evaluates the damage produced by the pipeline rupture which occurred in Edison, New Jersey, on March 23, 1994.

### 5.1.1 PAR 95-01

This Pipeline Accident Report (3) describes the rupture of a 36 inch natural gas transmission pipeline owned and operated by the Texas Eastern Transmission

Corporation. The rupture occurred at approximately 11:55 p.m. on March 23, 1994, on the property of Quality Materials, Inc. in Edison, New Jersey. Ignition of the escaping

Figure 5.1 - Location of Edison, New Jersey Accident
gas occurred within 2 minutes after the rupture, producing flames 400 to 500 feet high. While no deaths were directly attributed to the accident, the rupture produced extensive damage including the destruction of several buildings of a nearby apartment complex. The total cost of the damage exceeded 25 million dollars. The NTSB determined that the probable cause of the rupture was mechanical damage to the exterior pipeline surface. The damage reduced the pipeline wall thickness and probably resulted in a crack that grew to a critical size.

The accident location is shown in a copy of the aerial photograph presented as
Figure 5.1. This photograph (without the " + " symbol) was obtained from the NTSB (C. Batten: personal communication, February 1996), and was taken on March 30, 1994. In the photograph, north is approximately located to the upper right, and the location of the rupture is identified with the " + " symbol. The Durham Woods Apartments (i.e., the apartment complex which received significant damage) are located in the lower right of the photograph. An enlarged portion of this photograph depicts the rupture location (without the " + " symbol) and the Durham Woods Apartments, and is presented in Appendix C as Figure C.1. Another figure is provided (see Figure C.2) which is based on a sketch obtained from the NTSB (C. Batten: personal communication, February 1996) identifying the various buildings in the apartment complex.

As indicated previously, the rupture occurred at approximately 11:55 p.m., with ignition of the gas less than two minutes later. According to the PAR, the Edison Township Fire Department arrived at the accident scene at approximately 12:02 a.m. At that time, building number 12 was "fully involved" in fire, and the three buildings adjacent to it were quickly becoming involved. These are buildings numbered 9, 10 and 11 as shown on Figure C.2. Based on a discussion with Edison Township (personal communication, February 1996), buildings numbered 11, 12 and 24 were involved in fire upon arrival of the Fire Department, and buildings numbered 7, 8, 16, 19, 23 and 25 would have burned if these structures were not wetted down (it should be noted that buildings numbered 20, 21 and 22 were also destroyed by the accident).

If the buildings were set back from the rupture at a distance beyond the location of buildings $9,10,11,12$ and 24 (i.e., those buildings which were apparently involved in fire after a few minutes exposure to thermal radiation), then this distance would have provided protection for a few minutes from spontaneous ignition. This distance can be estimated using Equation 4-12, and verified by comparing the calculated burn radius to the actual location of the damage. However, since an appropriate scale could not be obtained for Figure C. 1 (since a scale was not provided with Figure 5.1), a scale for this figure must first be developed.

Figure C. 3 is an illustration of a portion of the accident location, and is based on a diagram that was likewise obtained from the NTSB (C. Batten: personal communication, February 1996). Using the existing scale on this figure, a landmark is selected which is clearly defined in both Figure C. 3 and Figure C. 1 and which can be used for scale comparison. The landmark chosen is the north-south length of buildings numbered 21 , 22 and 23 (which have similar lengths). Using the scale on Figure C.3, the length of these buildings is approximately 158 feet. The length of the image of building number 23 on the author's original Figure C. 1 is 0.625 inches. Therefore, the scale for Figure C. 1 is 158 feet per 0.625 inches, or 253 feet per inch. The white bar at the bottom of Figure C. 1 represents a length of 253 feet.

In order to use Equation 4-12, an appropriate heat flux must first be selected. Since the concern is protecting structures from spontaneous ignition for several minutes, a heat flux of $9,985 \mathrm{Btu} / \mathrm{hr} \mathrm{ft}^{2}$ from Appendix A will be selected. Inserting this value for heat flux into Equation 4-12, as well as the pipeline diameter of 36 inches and the incident operating pressure of 984.7 psia , the burn radius becomes:

$$
\begin{gathered}
\mathrm{BR}=36\{[(4,036.82)(984.7) /(9,985)]-37.52\}^{1 / 2} \\
=684 \text { feet }
\end{gathered}
$$

On Figure C.1, the distance from the rupture (taken as the midpoint of the open space between the two exposed ends of pipeline, or the center of the " + " symbol on Figure 5.1) and the approximate midpoint of the footprint for building number 9 (i.e., the building farthest from the rupture that was becoming involved in fire when the Fire Department arrived) is 772 feet. The estimated burn radius differs from the actual burn radius by less
than 12 percent. It can therefore be seen that the predicted burn radius does in fact approximate the distance to those buildings which were involved in fire a few minutes after the rupture.

A burn radius can likewise be estimated for determining the distance beyond which buildings will not ignite at all. For the Edison accident, this distance would extend beyond the location of those buildings which were wetted down. In order to predict the distance with Equation 4-12, a heat flux of $3,962 \mathrm{Btu} / \mathrm{hr} \mathrm{ft}^{2}$ is selected From Appendix A. This is the heat flux at which piloted ignition of wood occurs, so below this heat flux wood is not expected to ignite. From Equation 4-12, the burn radius becomes:

$$
\begin{gathered}
\mathrm{BR}=36\{[(4,036.82)(984.7) /(3,962)]-37.52\}^{1 / 2} \\
=1,119 \text { feet }
\end{gathered}
$$

Using the same scale as before, the distance from the rupture point to the midpoint of building number 25 (i.e., the building farthest from the rupture that was wetted down) is 1,101 feet. The predicted burn radius is therefore very similar to the actual distance determined from Figure C. 1, with a difference of less than 2 percent.

The above analysis has shown that Equation 4-12 provides a good estimation of burn radii for the Edison, New Jersey accident. The next section presents an analysis of the damage associated with an accident that occurred near Lancaster, Kentucky.
facing 61

Figure 5.2 - Location of Lancaster, Kentucky Accident

### 5.1.2 PAR 87-01

The accident described in this PAR (2) involves the rupture of a 30 -inch diameter pipeline owned by the Texas Eastern Gas Pipeline Company. The rupture and fire, which occurred near Lancaster, Kentucky, injured eight persons and resulted in \$500,000 damage to property adjacent to the pipeline. A depiction of the accident scene (reproduced from the PAR) is presented in Figure 5.2. The NTSB identified the probable cause of the accident as being the failure of the Texas Eastern Gas Pipeline Company to perform a comprehensive investigation of detected corrosion damage and to repair this damage.

The PAR indicates that the rupture occurred at 2:05 a.m. Based on the transcript of a hearing conducted by the NTSB (52), emergency responders arrived at the scene at 2:07 a.m and found that the area beginning at the Fletcher's trailer and extending to the structures behind the Henderson's home was involved in fire (see Figure 5.2). Specifically, during the aforementioned hearing, Mr. Mike Murphy of the Garrod County Disaster and Emergency Services states that upon arrival at the accident scene:
"The area from approximately the Fletcher's House Trailer back through this area, around the cross - - we could see quite easily that the Maggard Home was involved, that the Henderson Home was involved, the barns back in this area. So, basically, a large ball of fire in the area that was mentioned, however, there was a tremendous amount of exposure type fire from the other buildings that we have talked about - the barn. There was a significant amount of grass in the area on fire."

In a manner similar to that used for the Edison, New Jersey accident, a burn radius can be estimated using a heat flux of $9,985 \mathrm{Btu} / \mathrm{hr}^{2} \mathrm{f}^{2}$, since structures were ignited a few minutes after the rupture. Inserting this value once again into Equation 4-12, with the pipeline diameter of 30 inches and the incident operating pressure of 1001.7 psia, the burn radius becomes:

$$
\begin{gathered}
\mathrm{BR}=30\{[(4,036.82)(1,001.7) /(9,985)]-37.52\}^{1 / 2} \\
=575 \text { feet }
\end{gathered}
$$

Although Figure 5.2 does not have a scale, a scale can be determined using the 276 foot distance identified on the figure which corresponds to the distance from the rupture to the Henderson's house. On the original figure in the PAR, the 276 foot distance corresponds to 1.675 inches. Therefore, the appropriate scale is 276 feet per 1.675 inches, or 165 feet per inch. The black line in the lower right of Figure 5.2 represents this 165 foot length. Using this scale, the distance from the rupture to the midpoint of the footprint of the burned barn (i.e., the structure farthest from the rupture which was presumably buming a few minutes after the rupture) is 585 feet. The difference in the calculated versus the actual burn radius is less than 2 percent. As with the New Jersey accident, the predicted burn radius corresponds closely to the actual damage that was observed.

The information obtained for this accident is also sufficient for comparing burn radii estimated with Equation 4-12 to distances resulting in injuries to people. Based on the information provided in the hearing transcript (52), Mr. and Mrs. Henderson experienced second and third degree burns while their 12 year old daughter suffered
second degree burns. The Hendersons escaped from the scene of the rupture by leaving their house through the back door, running through a field and "over a bank at the river cliff" where the heat flux was low enough to allow them to rest.

Predicting the burn radii corresponding to a specific level of injury is more complex than predicting the distance for structural damage. Unlike structures, people can escape from the scene of a rupture and consequently experience heat fluxes which decrease with time and distance. Furthermore, escape velocities (i.e., the speed at which people run from the rupture) can vary as well as the level of clothing which the affected people are wearing (32). Although the exact path that the Henderson's took to escape is unknown, and the level of clothing they wore was not available in the information that was reviewed, it is known that they experienced a specific level of injury (i.e. second and third degree burns).

Using the specifications of the ruptured pipeline, one can show that Equation 4-12 predicts the occurrence of second and third degree burns to persons fleeing directly away from the rupture. First, a representative escape velocity must be determined. Hockey and Rew (32) provide escape velocities of $1 \mathrm{~m} / \mathrm{s}(3.28 \mathrm{ft} / \mathrm{s}), 2.5 \mathrm{~m} / \mathrm{s}(8.2 \mathrm{ft} / \mathrm{s}), 4 \mathrm{~m} / \mathrm{s}(13.1 \mathrm{ft} / \mathrm{s})$ and $6 \mathrm{~m} / \mathrm{s}(19.69 \mathrm{ft} / \mathrm{s})$. A representative escape velocity is considered to be $2.5 \mathrm{~m} / \mathrm{s}(8.2$ $\mathrm{ft} / \mathrm{s}$ ), which would reasonably be expected since people are fleeing the fire in the early morning hours after being asleep. Furthermore, it is assumed that the house provides shelter from the thermal radiation for the time in which the residents react to the accident (i.e., several seconds). From Figure 5.2, if a person exits from the rear of the Henderson's house (and, from Figure 5.2, knowing that the house is approximately 29 feet wide), that
person will be approximately $276+29=305$ feet from the rupture when first exposed to thermal radiation outside the house. After the first second, the person will be $305+8.2=313.2$ feet from the rupture; after two seconds the separation distance increases to $305+2(8.2)=321.4$ feet, and so on. Equation $4-12$ can be rewritten as follows in order to express heat flux as a function of separation distance, pipeline diameter and incident operating pressure:

$$
\begin{equation*}
\mathrm{K}=4036.82(\mathrm{P}) /\left[\left(\mathrm{BR}^{2} / \mathrm{D}^{2}\right)+37.52\right] \tag{5-1}
\end{equation*}
$$

Through use of Equation 5-1, the relationship of incident heat flux to exposure time can be examined for this accident. Furthermore, the following equation from Appendix A can be used for determining the average heat flux which will result in severe blistering:

$$
\begin{equation*}
\mathrm{q}_{2}=50 / \mathrm{t}_{\mathrm{c}}^{0.71} \tag{5-2}
\end{equation*}
$$

where:
$\mathrm{q}_{2}=$ Average heat flux $\left(\mathrm{kW} / \mathrm{m}^{2}\right)$
$\mathrm{t}_{\mathrm{c}}=$ Duration of exposure (seconds)
Table 5.1 provides the incident heat flux experienced over the first thirty seconds by a person fleeing the Lancaster, Kentucky rupture.

## Table 5.1

Heat Flux Experienced at Lancaster, Kentucky Accident

| Time Seconds | Distance from Rupture, feet | Heat Flux - Btu/hr $\mathrm{ft}^{2}$ |  |
| :---: | :---: | :---: | :---: |
|  |  | Equation 5-1 | Equation 5-2* |
| 0 | 305 | 28,703 | - |
| 2 | 321.4 | 26,552 | 9,689 |
| 4 | 337.8 | 24,610 | 5,923 |
| 6 | 354.2 | 22,856 | 4,442 |
| 8 | 370.6 | 21,269 | 3,621 |
| 10 | 387.0 | 19,829 | 3,090 |
| 12 | 403.4 | 18,521 | 2,715 |
| 14 | 419.8 | 17,330 | 2,434 |
| 16 | 436.2 | 16,244 | 2,214 |
| 18 | 452.6 | 15,252 | 2,036 |
| 20 | 469.0 | 14,343 | 1,889 |
| 22 | 485.4 | 13,510 | 1,766 |
| 24 | 501.9 | 12,740 | 1,660 |
| 26 | 518.3 | 12,035 | 1,568 |
| 28 | 534.7 | 11,385 | 1,488 |
| 30 | 551.1 | 10,784 | 1,417 |

*Values converted from $\mathrm{kW} / \mathrm{m}^{2}$ using the relationship $1 \mathrm{Btu} / \mathrm{hr}^{\mathrm{ft}^{2}}=3.1546 \mathrm{~kW} / \mathrm{m}^{2}$

In Figure C. 4 the information from Table 5.1 is plotted on the chart of heat flux versus exposure time. It can be seen that at all times considered, the estimated heat flux (calculated from Equation 5-1) is greater than the average heat flux necessary to cause severe blistering. Therefore, Equation 5-1 correctly predicts that persons fleeing the rupture will suffer at least second degree burns. Furthermore, it was previously mentioned that a heat flux of $9,510 \mathrm{Btu} / \mathrm{hr} \mathrm{ft}^{2}$ will cause third degree burns in 30 seconds. Equation 5-1 predicts that, at an escape velocity of $8.2 \mathrm{ft} / \mathrm{s}$, third degree burns should occur since the estimated heat flux is larger than $9,510 \mathrm{Btu} / \mathrm{hr}_{\mathrm{ft}^{2}}$. As documented by the NTSB, Mr. and Mrs. Henderson suffered both second and third degree burns.

While the available damage information from other pipeline accidents is not as extensive as for the two pipeline accidents that have been considered, the following section compares the available burn radii information to the estimates produced by Equation 4-12.

### 5.1.3 Other Data

While the information that was reviewed for the following accidents is not as detailed as the information obtained for the two accidents described above, burn radii estimated with Equation 4-12 are compared to the approximate burn radii that are based on the available damage information from other accidents documented by the NTSB. It should be noted that the following specific accidents were selected for comparison since a sufficient amount information is available to approximate burn radii and the incident operating pressures and pipeline diameters are within the ranges of Equation 4-12 (i.e., 575 10 1,200 psia and 14 inches to 36 inches in diameter).

1. PAR 71-01 (51) - Houston, Texas - September 9, 1969: This accident involved the rupture of a 14 inch pipeline at a pressure of approximately 789 psig ( 803.7 psia). Figure 2 in the PAR indicates that the roof of a house 300 feet from the rupture was scorched. In a manner similar to the analysis performed for the Edison, New Jersey accident, a heat flux of $3,962 \mathrm{Btu} / \mathrm{hr} \mathrm{ft}^{2}$ can be used to estimate the distance beyond which structures will not ignite. Using Equation 4-12, the burn radius becomes:

$$
\begin{gathered}
\mathrm{BR}=14\{[(4,036.82)(803.7) /(3,962)]-37.52\}^{1 / 2} \\
=391 \text { feet }
\end{gathered}
$$

The estimated burn radius represents a difference of 30 percent from the actual value.
2. PAR 75-02 (53) - Bealeton, Virginia - June 9, 1974: The accident, which involved the rupture of a 30 inch pipeline operating at 718 psig ( 732.7 psia ), produced a burn area 700 feet long and 400 feet wide. Figure 4 of the PAR indicates that the burn area was not symmetrical, but rather downstream of the rupture. Therefore, the actual burn radius will be considered to be 700 feet. The estimated burn radius is found as follows, using Equation 4-12 with a heat flux of $3,962 \mathrm{Btu} / \mathrm{hr}^{2}{ }^{2}$ (i.e. the heat flux for piloted ignition of wood, which is assumed to be the heat flux associated with the extent of a burn area):

$$
\begin{gathered}
\mathrm{BR}=30\{[(4,036.82)(732.7) /(3,962)]-37.52\}^{1 / 2} \\
=799 \text { feet }
\end{gathered}
$$

This value represents a difference of 14 percent from the actual value.
3. PAR 77-01 (50) - Cartwright, Louisianna - August 9, 1976: The rupture of a 20 -inch pipe operating at 770 psig ( 784.7 psia ) resulted in a burn area of approximately 12 acres. It is unclear from the information that was reviewed whether or not the burn area was symmetric. Figure 2 of the PAR indicates damage surrounding the pipeline. Complicating matters is the fact that horizontal flames vented from the rupture for a distance of 100 to 150 feet. If a symmetric burn area is considered, then a burn radius can be found by first calculating the
total burn area in square feet. Since one acre represents 43,560 square feet, the total burn area is $12(43,560)=522,720 \mathrm{ft}^{2}$. From the equation for the area of a circle (i.e., area $=\pi[\text { radius }]^{2}$ ), the actual burn radius becomes:

$$
\mathrm{BR}=(522,720 / \pi)^{1 / 2}=408 \text { feet }
$$

By comparison, use of Equation $4-12$ with a heat flux of $3962 \mathrm{Btu} / \mathrm{hr} \mathrm{ft}^{2}$ produces the following:

$$
\begin{gathered}
\mathrm{BR}=20\{[(4,036.82)(784.7) /(3,962)]-37.52\}^{1 / 2} \\
=552 \text { feet }
\end{gathered}
$$

For this accident, the estimated burn radius differs from the actual value by 35 percent.
4. DCA-85-FP-003 (54) - Jacksonville, Louisianna - November 25, 1984: The pipeline involved in this accident had a diameter of 30 inches and was operating at $1,016 \mathrm{psig}(1,030.7 \mathrm{psia})$. A non-symmetrical damage area was produced, with the rupture incinerating an area 900 feet north, 500 feet south and 180 feet to the east and west of the rupture. If the burn radius is considered to be the maximum linear distance from the rupture to the edge of the incinerated area, the radius is then 900 feet. Again, using Equation 4-12 and a heat flux of $3,962 \mathrm{Btu} / \mathrm{hr}_{\mathrm{ft}}{ }^{2}$, the estimated burn radius becomes:

$$
\begin{gathered}
\mathrm{BR}=30\{[(4,036.82)(1030.7) /(3,962)]-37.52\}^{1 / 2} \\
=955 \text { feet }
\end{gathered}
$$

The estimated burn radius differs from the actual burn radius by only 6 percent.
5. DCA-86-FP-009 (49) - Cale, Arkansas - February 24, 1986: The information available for this accident indicates that a 20 inch pipeline operating at 750 psig ( 764.7 psia ) ruptured and resulted in a burn area of 15 acres. Since 1 acre is equivalent to $43,560 \mathrm{ft}^{2}$, the burn area (in square feet) is $15(43,560)=653,400 \mathrm{ft}^{2}$. Since details of the actual burn shape were not available, a circular burn patten is considered, and the actual burn radius becomes (using the equation for the area of a circle):

$$
\mathrm{BR}=(653,400 / \pi)^{1 / 2}=456 \text { feet }
$$

Using a heat flux of $3,962 \mathrm{Btw} / \mathrm{hr}^{\mathrm{ft}}{ }^{2}$, Equation 4-12 estimates the burn radius to be:

$$
\begin{gathered}
\mathrm{BR}=20\{[(4,036.82)(764.7) /(3,962)]-37.52\}^{1 / 2} \\
=545 \text { feet }
\end{gathered}
$$

The difference in actual versus estimated values is approximately 20 percent.
6. PAR 87-01 (2) - Beaumont, Kentucky - April 27, 1985: This accident (described in the same report as the Lancaster, Kentucky accident) involved the rupture of a 30 inch diameter pipeline operating at 990 psig ( $1,004.7$ psia). The accident incinerated an area 700 feet long and 500 feet wide. From Figure 3 in the PAR, the burn area is not truly symmetric; rather, the longest dimension from the rupture is approximately 500 feet. If this dimension is considered the actual burn radius, then the estimated burn radius is:

$$
\begin{gathered}
\mathrm{BR}=30\{[(4,036.82)(1,004.7) /(3,962)]-37.52\}^{1 / 2} \\
=942 \text { feet. }
\end{gathered}
$$

This estimation represents the largest difference observed between predicted and actual burn radii, with a value of 88 percent.

### 5.2 Comparison to Separation Distances Developed through Hazard Analysis

### 5.2.1 Separation Distances Imposed in the Dutch Regulations

Earlier chapters discussed the Dutch regulations and the two types of separation distances that have been developed. The following discussion is based on personal communications from Hans van Poelje of N.V. Nederlandse Gasunie dated November 30 , 1995; June 21, 1996; and September 30, 1996. The first type of separation distance is called a proximity, or building distance. This is the distance between a pipeline and residential buildings (or special structures) and corresponds to a $10^{-6}$ individual risk. The second type of distance is called a survey, or effect distance. This distance is determined for the purpose of identifying the location classification and corresponds to a $10^{-8}$ individual risk.

Based upon the available information, the building distances imposed in the Dutch regulations can be compared to distances that are estimated through use of Equation 4-12. However, since Equation 4-12 was not developed using risk assessment
techniques, a basis for comparison must first be selected. For comparison to building distances, Equation 4-12 will use as a basis a heat flux of $9,985 \mathrm{Btu} / \mathrm{hr}_{\mathrm{ft}}{ }^{2}$. As was shown earlier in this chapter, this level of thermal radiation results in the destruction of buildings after a few minutes exposure.

The Dutch regulations specify three ranges of operating pressures (20 to 50 bar; 50 to 80 bar; and 80 to 110 bar) and pipeline diameters from 2 inches to 48 inches. In order to provide a meaningful comparison, the range of pipeline diameters and pressures used in the development of Equation 4-12 will be evaluated (i.e., diameters from 14 to 36 inches and pressures from 575 to 1,200 psia). By converting the Dutch pressure ranges from bar to psia (by using the relationship 1.01325 bar equals 14.696 psia, and converting gauge to absolute pressure), the pressure ranges become: 304.8 to 739.9 psia; 739.9 to $1,175.0$ psia; and $1,175.0$ to $1,610.1$ psia. Although the midpoints of these ranges are approximately 522.3 psia, 957.4 psia, and $1,392.6 \mathrm{psia}$, the three pressures which will be used in Equation 4-12 are 522.3 psia; 957.4 psia and $1,200.0$ psia. The pressure of $1,200.0 \mathrm{psia}$ is used in lieu of $1,392.6$ psia since $1,200.0$ psia represents the upper limit of pressure used to develop Equation 4-12, yet still lies within the third range.

The comparisons are presented in Table C. 1 of Appendix C. It can be seen that Equation 4-12 estimates significantly larger separation distances than the building distances determined by the Dutch. However, as shown through the analyses of the

Edison, New Jersey and Lancaster, Kentucky accidents, the building distances developed by the Dutch will not be protective of structures. For a 36 -inch diameter pipeline, the maximum building distance is 148 feet. This distance would clearly not have been protective of structures for the two aforementioned accidents.

Although the building distances and burn radii do not correspond, the trends in both at conditions of constant pressure (with varying diameter) and constant diameter (with varying pressure) do correspond closely. If pressure is held constant, then the following ratio is produced when using Equation 4-12:

$$
\begin{gather*}
\left.\left.\mathrm{BR}_{1} / \mathrm{BR}_{2}=\left[\mathrm{D}_{1}\{(4,036.82) \mathrm{P} / \mathrm{K})-37.52\right\}^{1 / 2}\right] /\left[\mathrm{D}_{2}\{(4,036.82) \mathrm{P} / \mathrm{K})-37.52\right\}^{1 / 2}\right] \\
=\mathrm{D}_{1} / \mathrm{D}_{2} \tag{5-3}
\end{gather*}
$$

If diameters are held constant, then Equation 4-12 produces the following ratios:

$$
\begin{align*}
\mathrm{BR}_{1} / \mathrm{BR}_{2} & \left.\left.=\left[\mathrm{D}\left\{(4,036.82) \mathrm{P}_{1} / \mathrm{K}\right)-37.52\right\}^{1 / 2}\right] /\left[\mathrm{D}\left\{(4,036.82) \mathrm{P}_{2} / \mathrm{K}\right)-37.52\right\}^{1 / 2}\right] \\
& =\left[\left\{\left((4,036.82) \mathrm{P}_{1} / \mathrm{K}\right)-37.52\right\} /\left\{\left((4,036.82) \mathrm{P}_{2} / \mathrm{K}\right)-37.52\right\}\right]^{1 / 2} \tag{5-4}
\end{align*}
$$

Tables C. 2 and C. 3 of Appendix C present the comparison of building distances for constant pressure and diameter. It can be seen that both the Dutch approach and Equation 4-12 produce very similar trends (i.e. similar ratios) whether conditions of constant pressure or constant diameter are evaluated.

### 5.2.2 Separation Distances Determined by the British Health and Safety Executive

Hill and Catmur performed a study for the British Health and Safety Executive (31) to evaluate how risks from various hazardous pipelines compared. As part of the above study, distances from a vertical flame jet to a heat flux level of $10 \mathrm{~kW} / \mathrm{m}^{2}$ $\left(3,170 \mathrm{Btu} / \mathrm{hr} \mathrm{ft}^{2}\right)$ are provided for the pipelines under consideration. The authors simulated the flame jet as an inclined line source with a receptor 1.5 meters ( 4.92 feet) high at ground level. Furthermore, the authors indicate that the model which was used provides the maximum view factor between the source and receptor, with the thermal radiation being a function of the flame's emissivity, the transmissivity of the air, the view factor and the radiant energy of the burning compound.

A comparison can be made between those distances and the distances estimated using Equation 4-12. This comparison is presented in Table 5.2, for pipelines with diameters within the range of 14 inches to 36 inches, and operating pressures within the range from 500 psia to $1,200 \mathrm{psia}$ :

## Table 5.2

Comparison of Health and Safety Executive Results with Equation 4-12

| Pipeline |  |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: |
| Diameter | Pressure |  | Separation Distances-Feet |  | Percent |
| Inches | Barg | Psia | British* | Equation 4-12 | Difference |
| 24 | 70 | 1030.0 | 820 | 857 | 4.5 |
| 16 | 70 | 1030.0 | 564 | 571 | 1.2 |

[^0]The separation distances found from Equation 4-12 agree very closely with those predicted by Hill and Catmur. If the ranges of diameters and pressures is extended to include all of the natural gas pipelines in the study, the comparison in Table 5.3 is obtained:

## Table 5.3

Comparison of all Natural Gas Pipelines

| Pipeline |  |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: |
| Diameter | Pressure |  | Separation Distances-Feet |  | Percent |
| Inches | Barg | $\underline{\text { Psia }}$ | British ${ }^{*}$ | Equation 4-12 | Difference |
| 42 | 70 | 1030.0 | 1,385 | 1,499 | 8.2 |
| 24 | 70 | 1030.0 | 820 | 857 | 4.5 |
| 16 | 70 | 1030.0 | 564 | 571 | 1.2 |
| 6 | 70 | 1030.0 | 226 | 214 | 5.3 |
| 24 | 16 | 246.8 | 443 | 399 | 9.9 |
| 24 | 7 | 116.2 | 351 | 252 | 28.2 |
| 6 | 16 | 246.8 | 138 | 100 | 27.5 |
| 6 | 7 | 116.2 | 95 | 63 | 33.7 |

It can be seen that even outside the range of diameters and pressures for which Equation 4-12 was developed that this relationship still produces results which approximate those of the British. The higher percent differences reflected in the last three entries of Table 5.3 are probably due to the use of low operating pressures, either singly or in combination with small diameters, which are outside the range for which Equation 4-12 was developed.

In summary, the analyses performed in this chapter have shown that the methodology developed in this thesis can be used to provide estimations of burn radii which are consistent with the burn radii determined for those accidents reported upon by the NTSB where there is ample documentation to conduct an analysis using the approach developed herein. Additionally, this methodology produces results consistent with the work published by the British Health and Safety Executive. While the burn radii estimated with Equation 4-12 are not similar to the building distances developed by the Dutch, it is clear from analyses performed in this chapter that the distances developed by the Dutch would not be protective of buildings. However, the trends in both the building distances and the distances estimated with Equation 4-12 (i.e., the variation of distance with both pressure and diameter) are found to be very consistent.

The next chapter provides additional evaluation of the methodology developed in this thesis. Specifically, an analysis will be performed in which the sensitivity of the burn radius is evaluated based on variations of the assumptions stated in Chapter 4.

## CHAPTER 6

## RESULTS AND DISCUSSION

The previous chapters have described the work that was performed for this thesis. This work included the development and testing of a relationship for the estimation of burn radii based on a pipeline's diameter and incident operating pressure. Consequently, the primary result of this thesis is Equation 4-12:

$$
\begin{equation*}
\mathrm{BR}=\mathrm{D}\{(4,036.82) \mathrm{P} / \mathrm{K})-37.52\}^{1 / 2} \tag{4-12}
\end{equation*}
$$

where
$\mathrm{BR}=$ Burn Radius ( ft )
$\mathrm{D}=$ Pipeline Diameter (in)
$\mathrm{P}=$ Incident operating pressure (psia)
$\mathrm{K}=$ Heat flux (Btu/hr $\mathrm{ft}^{2}$ )
Through the use of Equation 4-12, charts have been developed which can be used to estimate burn radii. These charts are provided in Appendix B for heat flux values (listed in Appendix A) of $3,962 \mathrm{Btu} / \mathrm{hr}_{\mathrm{ft}}{ }^{2}$ (piloted ignition of wood); $6,340 \mathrm{Btu} / \mathrm{hr}_{\mathrm{ft}}{ }^{2}$ (blistering of bare skin in 4 seconds and $1 \%$ lethality in 20 seconds); 9,510 Btu/hr $\mathrm{ft}^{2}$ (causes third degree burns in 30 seconds); and 9,985 Btu/hr $\mathrm{ft}^{2}$ (spontaneous ignition of wooden structures after a few minutes). Equation 4-12 and the charts in Appendix B are applicable to natural gas pipelines with diameters ranging from 14 inches to 36 inches and incident operating pressures in the range of 575 psia to $1,200 \mathrm{psia}$. The above ranges of diameters and pressures are applicable to most major natural gas transmission pipelines operating in the United States.

In Chapter 5, the sensitivity of Equation 4-12 to changes in pressure (at constant diameter) and to changes in diameter (at constant pressure) was discussed in relation to the work performed by the Dutch. In this chapter, an analysis will be performed to evaluate the sensitivity of Equation 4-12 to variations in heat flux. Furthermore, this analysis will be expanded to include the effects of variations in atmospheric transmissivity, wind direction and wind speed on the estimation of the burn radius.

### 6.1 Effect of Variation in Heat Flux on Burn Radius

Selection of an appropriate heat flux for use in Equation 4-12 is critical in the estimation of an appropriate burn radius. It is therefore worthwhile to evaluate the effect that a variation in heat flux can have on the burn radius. Equation 4-12 was developed for use with natural gas pipelines having diameters from 14 inches to 36 inches and incident operating pressures from 575 psia to 1,200 psia; therefore, the following conditions (which span the range of diameters and pressures) will be used in the analysis:

## Table 6.1

Pipeline Conditions

| Pipeline Diameter | Incident Operating Pressure |
| :--- | :--- |
|  |  |
| 14 Inches | 575 psia |
| 14 Inches | $1,200 \mathrm{psia}$ |
| 36 Inches | 575 psia |
| 36 Inches | $1,200 \mathrm{psia}$ |

In addition, since heat flux values from $3,962 \mathrm{Btu} / \mathrm{hr}_{\mathrm{ft}^{2}}$ to $9,985 \mathrm{Btu} / \mathrm{hr} \mathrm{ft}^{2}$ were shown in Chapter 5 to provide results (i.e., burn radii estimates) similar to those observed during actual pipeline ruptures, the analysis will be limited to values within this range.

Through use of Equation 4-12, burn radii are estimated for each of the pipeline conditions specified in Table 6.1, at the following heat flux levels: $3,962 \mathrm{Bu} / \mathrm{hr}^{\mathrm{ft}}{ }^{2} ; 6,340$ $\mathrm{Btu} / \mathrm{hr} \mathrm{ft}^{2} ; 9,510 \mathrm{Btu} / \mathrm{hr}^{2}{ }^{2}$; and 9,985 Btu/hr $\mathrm{ft}^{2}$. Table 6.2 provides a tabular representation of this information.

Table 6.2
Burn Radii at Various Heat Fluxes

| Heat Flux | Pipeline Conditions - Diameter/Incident Operating Pressure |  |  |  |
| :---: | :---: | :---: | :---: | :---: |
| Btu/hr $\mathrm{ft}^{2}$ | 14"/575psia | 14"/1,200 psia | 36'/575psia | 36"/1,200psia |
| 3,962 | 328 | 482 | 843 | 1,239 |
| 6,340 | 254 | 377 | 653 | 970 |
| 9,510 | 201 | 304 | 517 | 782 |
| 9,985 | 195 | 296 | 503 | 762 |

From the above results, it can be seen that the burn radius varies from a minimum of 195 feet to a maximum of 1,239 feet. The variation of burn radius with heat flux is evident from the columns of Table 6.2 , which illustrate the decrease in burn radius with increasing heat flux for each of the four pipeline conditions. The use of Equation 4-12 necessitates the selection of a single value of heat flux, which is presumed to be associated with a specific level of thermal radiation damage. However, from the heat flux information presented in Appendix A, it can be seen that a range of heat flux values could be associated with a particular level of damage. Therefore, it is appropriate to examine
the effect that a variation in heat flux has on the estimated burn radius. Table 6.3 provides the results of burn radius calculations for each of the four heat flux values of Table 6.2, in which each of the heat flux values is varied (i.e., increased and decreased) by 25 percent. This amount of variation is consistent with the variation of heat flux values in Appendix A, for a specified level of thermal radiation damage.

Table 6.3
Variation of Heat Flux Values
Burn Radius and Percent Difference* for

| Heat Flux | Pipeline Conditions - Diameter/Incident Operating Pressure |  |  |  |
| :---: | :---: | :---: | :---: | :---: |
| Btu/hr $\mathrm{ft}^{2}$ | 14"/575psia | 14"/1.200 psia | 36"/575 psia | $36^{\prime \prime} / 1.200 \mathrm{psia}$ |
| 2,972 | 381(16) | 559(16) | 982(16) | 1,437(16) |
| 3,962 | 328 | 482 | 843 | 1,239 |
| 4,953 | 291(-11) | 429(-11) | 747(-11) | 1,104(-11) |
| 4,755 | 297(17) | 439(16) | 764(17) | 1,128(16) |
| 6,340 | 254 | 377 | 653 | 970 |
| 7,925 | 224(-12) | 335(-11) | 575(-12) | 862(-11) |
| 7,133 | 238(18) | 355(17) | $611(18)$ | 912(17) |
| 9,510 | 201 | 304 | 517 | 782 |
| 11,888 | 176(-12) | 269(-12) | 452(-13) | 692(-12) |
| 7,489 | 231(18) | 346(17) | 594(18) | 889(17) |
| 9,985 | 195 | 296 | 503 | 762 |
| 12,481 | 171(-12) | 262(-11) | 439(-13) | 674(-12) |

*Burn radius is the first value specified in the table; percent difference is parenthesized.

From Table 6.3, it can be seen that over and under estimating the heat flux by 25
percent can result in maximum variations in the burn radius of -13 percent and 18 percent, respectively. These variations are considered reasonable in light of the large numerical variation in the heat flux values (i.e., variations of at least $990 \mathrm{Btu} / \mathrm{hr}_{\mathrm{ft}}{ }^{2}$ ).

In Chapter 4, it was explained that Equation 4-12 does not account for the heat flux due to solar radiation. Appendix A indicates that this heat flux varies from 250 to $380 \mathrm{Btu} / \mathrm{hr} \mathrm{ft}^{2}$. If a representative value for solar heat flux is selected to be $315 \mathrm{Btu} / \mathrm{hr}_{\mathrm{ft}^{2}}$ (i.e., the midpoint of the solar heat flux range), then a comparison of the burn radii with and without solar radiation can be made. This comparison is shown in Table 6.4, in which the burn radii for the four pipeline conditions are calculated with and without the contribution of the solar heat flux.

Table 6.4
Variation of Heat Flux Values (Contribution of Solar Heat Flux)

|  | $\begin{array}{c}\text { Burn Radius and Percent Difference* for }\end{array}$ |  |  |  |  |
| :--- | :--- | :--- | :--- | :--- | :---: |
| $\begin{array}{llll}\text { Heat Flux } \\ \text { Btu/hr } \mathrm{ft}^{2}\end{array}$ | Pipeline Conditions - Diameter/Incident Operating Pressure |  |  |  |  |$)$

*Burn radius is the first value specified in the table; percent difference is parenthesized.

The results in Table 6.4 indicate that a variation in heat flux of several hundred Btu's/hr $\mathrm{ft}^{2}$ will not have a significant impact on the burn radius. Therefore, the assumption to exclude solar radiation during the development of Equation 4-12 is appropriate.

### 6.2 Effect of Variation in Atmospheric Transmissivity on Burn Radius

During development of Equation 4-12, the following relationship was used to provide an estimate of atmospheric transmissivity:

$$
\begin{equation*}
\tau=0.79(100 / \mathrm{r})^{1 / 16}(100 / \mathrm{D})^{1 / 16} \tag{4-4}
\end{equation*}
$$

where
$r=$ relative humidity, $\%$
$\mathrm{D}=$ distance to flame, feet

As explained in Chapter 4, the authors of this equation (36) caution that Equation $4-4$ is "strictly applicable" when the flame is radiating at a temperature of $2240^{\circ} \mathrm{F}$, relative humidity is greater than $10 \%$, dry bulb temperature is $80^{\circ} \mathrm{F}$ and the distance from the flame is greater than 100 feet but less than 500 feet. However, the authors also indicate that an "order of magnitude" estimate can be made under a wider range of conditions. For purposes of simplifying Equation 4-12, a value for atmospheric transmissivity was determined assuming a distance from the flame of 500 feet and a relative humidity of 50 percent. Using Equation 4-4, this value was calculated to be 0.746.

The sensitivity of the burn radius to changes in atmospheric transmissivity can be evaluated at the four conditions identified in Section 6.1 by varying both the relative humidity and the distance used in Equation 4-4. Table 6.5 provides the atmospheric transmissivities that are obtained from Equation 4-4 when relative humidities range from 10 percent to 100 percent and distances vary from 100 to 500 feet:

Table 6.5
Atmospheric Transmissivities

| Humidity, <br> Percent | $\underline{y}$ | Distance from Flame - Feet <br> 200 |  |  |  |
| :--- | :--- | :--- | :--- | :--- | :--- |
| $l l l l l l$ |  |  |  |  |  |
| 10 | 0.912 |  | $\underline{300}$ | $\underline{400}$ | $\underline{500}$ |
| 20 | 0.874 | 0.874 | 0.837 | 0.816 | 0.837 |
| 30 | 0.852 | 0.816 | 0.795 | 0.781 | 0.825 |
| 40 | 0.837 | 0.801 | 0.781 | 0.767 | 0.790 |
| 50 | 0.825 | 0.790 | 0.770 | 0.757 | 0.757 |
| 60 | 0.816 | 0.781 | 0.762 | 0.748 | 0.746 |
| 70 | 0.808 | 0.774 | 0.754 | 0.741 | 0.731 |
| 80 | 0.801 | 0.767 | 0.748 | 0.735 | 0.724 |
| 90 | 0.795 | 0.762 | 0.742 | 0.729 | 0.719 |
| 100 | 0.790 | 0.757 | 0.738 | 0.724 | 0.714 |

Table 6.5 indicates that for relative humidities from 10 to $100 \%$ and distances from 100 to 500 feet, the transmissivity varies from 0.912 to 0.714 . It can be seen that the atmospheric transmissivity decreases with increasing relative humidity and distance, since for these conditions a larger amount of thermal radiation is absorbed by the atmosphere.

It was previously mentioned that the burn radius produced by Equation 4-12 varies from 195 feet to 1,239 feet. The range of possible atmospheric transmissivities associated with Equation 4-12 can be found by determining the transmissivities at the smallest burn radius (and lowest applicable relative humidity) and at the highest burn radius (and highest applicable relative humidity). The smallest burn radius and applicable relative humidity is respectively considered to be 195 feet and 10 percent. Using Equation 4-4, the atmospheric transmissivity is:

$$
\begin{gathered}
\tau=0.79(100 / 10)^{1 / 16}(100 / 195)^{1 / 16} \\
=0.875
\end{gathered}
$$

Since the maximum burn radius of 1,239 feet is outside the range for which Equation 4-4 is strictly applicable, (i.e., 100 to 500 feet), the value of atmospheric transmissivity at this distance is estimated using a general plot of transmissivity versus path length (20). As indicated in (20), this plot is for a source temperature of 1400 K . which is the general flame temperature for the burning of hydrocarbon fuels. From the plot, the value of $\tau$ for a path length of 1,239 feet (or 378 meters) using the dashed line representing 100 percent relative humidity (which will result in a minimum value of $\tau$ ) is approximately 0.470 . Therefore, the atmospheric transmissivity can vary from a maximum of 0.875 to a minimum of 0.470 . Equation $4-12$ can be rewritten to make $\tau$ a variable by dividing the $(4,036.82) \mathrm{P} / \mathrm{K}$ term by 0.746 (the previously assumed value of $\tau$ ); the equation then becomes:

$$
\begin{align*}
\mathrm{BR} & =\mathrm{D}\{(4,036.82 / 0.746) \mathrm{P} / \mathrm{K})-37.52\}^{1 / 2} \\
& =\mathrm{D}\{(5,411.29 \tau) \mathrm{P} / \mathrm{K})-37.52\}^{1 / 2} \tag{6-1}
\end{align*}
$$

The burn radius can now be computed at the four pipeline conditions using the maximum, minimum and previously assumed values of $\tau$. The results of these calculations are provided in Table D. 1 of Appendix D. In addition, Table D. 2 of Appendix D provides the percent difference between the burn radii that are calculated using the assumed value of $\tau(0.746)$ and the burn radii that are calculated using the maximum and minimum values of $\tau(0.875$ and 0.470 , respectively). From Tables D. 1 and D.2, it can be seen that Equation 4-12 can under predict and over predict the burn radius by 9 percent and 34 percent, respectively. While an overestimation of the burn radius by approximately a third may initially appear to be significant, it must be remembered that this percentage corresponds to extreme conditions (i.e., relative humidity of 100 percent and a distance of 1,239 feet from the flame). Furthermore, since these percentages correspond to a range, the actual percent difference in any situation would be expected to lie within this range and be less than the extreme values.

The combined effects of variations in heat flux and atmospheric transmissivity can be shown through calculation of burn radii at the four pipeline conditions and heat flux values identified in Table 6.3, using atmospheric transmissivities which vary from 0.470 to 0.875 . Table D. 3 in Appendix D provides the calculated burn radii at all of the aforementioned conditions, while Table D. 4 lists the percentage differences of these variations from burn radii calculated at the previously described conditions (i.e., pipeline diameters of 14 inches and 36 inches; incident operating pressures of 575 psia and 1,200
psia; heat flux values of $3,962 \mathrm{Btu} / \mathrm{hr} \mathrm{ft}^{2} ; 6,340 \mathrm{Btu} / \mathrm{hr} \mathrm{ft}^{2 ;} 9,510 \mathrm{Btu} / \mathrm{hr} \mathrm{ft}^{2 ;}$ and 9,985 $\mathrm{Btu} / \mathrm{hr} \mathrm{ft}^{2}$; and a value of 0.746 for atmospheric transmissivity). The largest variations are seen for the combination of low $\tau$ and high $K$ (i.e., "low" and "high" in relation to the pipeline condition being considered); this combination results in underprediction of the burn radius by an average of 33 percent. However, a combination of high $\tau$ and high K result in an underprediction which averages only 3 percent. Conversely, a combination of high $\tau$ and low K produces overpredictions of burn radii averaging 27 percent, while combining a low $\tau$ with a low K results in an average underprediction of 9 percent. It must again be noted that these percentages correspond to extreme conditions, and that the actual percent difference in any situation would lie within ranges bounded by these values.

### 6.3 Effect of Variation in Wind Speed and Direction on Burn Radius

As described previously, Equation 4-12 is developed assuming a vertical flame (i.e., no wind) and, consequently, a symmetrical burn area. Graphical information has been obtained from the Canadian TSB (L. Gales: personal communication, September 10, 1996) which relates heat flux to distance at wind speeds of 10 kilometers per hour ( 6.2 miles per hour) and 20 kilometers per hour ( 12.4 miles per hour) at upwind, downwind and crosswind locations. The wind speeds cited herein are appropriate for this discussion, since they span a range which is representative of many cities in the United States. As an example, Table 6.6 provides the mean yearly wind speeds for several major cities (55).

## Table 6.6

Wind Speeds for Several U.S. Cities

| City | Mean Yearly Wind Speed, <br> miles per hour |
| :--- | :--- |
| Atlantic City, NI | 10.1 |
| Newark, NJ | 10.2 |
| Philadelphia, PA | 9.5 |
| Charleston, SC | 8.6 |
| Madison, WI | 9.8 |
| Nashville, TN | 8.0 |
| Columbus, OH | 8.5 |
| Helena, MT | 7.7 |
| Shreveport, LA | 8.4 |
| Orlando, FL | 8.5 |

The graphs were prepared for pipeline diameters of 24,36 and 42 inches at a pressure of $1,000 \mathrm{psig}(1014.7 \mathrm{psia})$ and, with the exception of graphs for the 42 inch diameter pipeline (which is outside the diameter range of Equation 4-12), have been provided in Appendix E. It should be noted that the graphs show the progression of heat flux with time, with the curve for time zero at the top. It is this curve which produces the maximum heat flux at a specified distance.

Burn radii at heat flux values of $9,510 \mathrm{Btu} / \mathrm{hr} \mathrm{ft}^{2} ; 7,925 \mathrm{Btu} / \mathrm{hr} \mathrm{ft}^{2} ; 6,340$ $\mathrm{Btu} / \mathrm{hr} \mathrm{ft}^{2}$ and $4,755 \mathrm{Btu} / \mathrm{hr}_{\mathrm{ft}^{2}}$ were obtained from the time zero curves in Appendix E and compared to distances determined using Equation 4-12. These heat flux values were selected since they lie within the previously identified range of $3,962 \mathrm{Btu} / \mathrm{hr}^{2} \mathrm{ft}^{2}$ to 9,985 Btu/hr $\mathrm{ft}^{2}$ for which Equation 4-12 is applicable. Furthermore, these heat flux values facilitate reading distances from the figures in Appendix E.

The comparison of distances and the calculation of percentage differences is respectively provided in Tables D. 5 and D. 6 (for the 24 inch diameter pipeline) and Tables D. 7 and D. 8 (for the 36 inch diameter pipeline). The distances listed in Tables D. 5 and D. 7 show that the burn radius is not symmetrical when there is wind, and that the radius is greater in the downwind direction than in the upwind direction. From Tables D. 6 and D. 8 , it can be seen that the largest percentage difference between the Canadian distances and the distances estimated with Equation 4-12 occurs for the upwind condition ( 56 percent for the 36 inch diameter pipeline and 32 percent for the 24 inch diameter pipeline), and that for upwind conditions the percentages decrease with decreasing wind speed. However, an assumption was previously made that Equation 4-12 will predict the longest dimension of a burn area, which would be the distance associated with the downwind condition. The maximum percent difference between the Canadian distances and the distances estimated using Equation $4-12$ is less if the downwind condition is considered: 27 percent for the 36 inch pipeline and -16 percent for the 24 inch pipeline. Furthermore, for all locations and wind speeds the percent difference decreases with decreasing heat flux (so that overestimations are associated with a heat flux of $9,510 \mathrm{Btu} / \mathrm{hr}_{\mathrm{ft}}{ }^{2}$ and underestimations can appear when considering lower values of heat flux). However, given the variability in weather conditions which may be present during an actual pipeline rupture, the differences in burn radius cited herein are considered to be reasonable.

In summary, the analyses performed in this chapter serve to illustrate that the assumptions made during the development of Equation 4-12 are appropriate. Equation 4-12 and the charts of Appendix $B$ can be used to provide estimations of burn radii for various pipeline diameters and incident operating pressures. However, it must be stressed that safe separation distances determined through the use of Equation 4-12 or the Appendix B charts are estimations. There are numerous variables, several of which have been considered in this chapter, which will influence the burn radius associated with a pipeline rupture. The advantage to using the method described in this thesis is that the method is straightforward and provides reasonable estimates of actual burn radii.

The next chapter provides a summary of this thesis through several concluding remarks. Recommendations are also made for additional work, which can be used to refine the procedure for estimating safe separation distances.

## CHAPTER 7

## CONCLUSIONS AND RECOMMENDATIONS

The work described in this thesis has led to the development of a method for estimating safe separation distances from natural gas transmission pipelines. This method was verified based on information from actual pipeline accidents, and provides a means to determine the safe separation distance, as a measurement of the burn radius, through knowledge of the pipeline's diameter and incident operating pressure. The method can be used by regulators to determine the distances at which future development might be placed from existing pipelines or, perhaps more realistically, the method can be used to evaluate appropriate incident operating pressures for pipelines which traverse densely populated areas.

The procedure described in this thesis is easy to apply and does not require extensive computational efforts. The basis for estimation is Equation 4-12, which describes the burn radius as a function of a pipelines's diameter and incident operating pressure. As illustrated in the charts found in Appendix B of this thesis, Equation 4-12 is readily represented in graphical form. The method is applicable to pipelines with diameters ranging from 14 inches to 36 inches and incident operating pressures from 575 psia to $1,200 \mathrm{psia}$, which constitute the mass majority of natural gas transmission pipelines in service in the United States. For levels of thermal radiation damage corresponding to heat flux values from $3,962 \mathrm{Btu} / \mathrm{hr} \mathrm{ft}^{2}$ to $9,985 \mathrm{Btu} / \mathrm{hr} \mathrm{ft}^{2}$, the method will predict safe separation distances ranging from 195 feet to 1,200 feet. The range of heat flux values noted above are applicable to the major consequences to life, limb and property that should be of interest to most analysts.

Several recommendations can be made for future research endeavors that can refine the work that has been performed herein. As advances are made in the areas of fire protection and hazard assessment, new information may be developed concerning the relationship between heat flux levels and corresponding levels of themal radiation damage. This new information can then be used to prepare burn radius charts for which the heat flux level corresponding to a specified level of damage is more accurately known. Another recommended area for future work includes the refinement of the procedure to predict atmospheric transmissivity, as the current method only provides order of magnitude estimates at distances beyond 500 feet.

A recommendation can also be made with regard to the reporting of pipeline accident data, since the reports produced by the NTSB and the Canadian TSB vary in the type and amount of accident-related information. Consistent reporting of data such as the volume of gas released, the duration of release and flame characteristics would be beneficial, in that the model developed in this thesis can be tested more accurately. Furthermore, this additional information may also be used to expand the diameter and pressure ranges for which Equation 4-12 is applicable

Although there are areas amenable to refinement, the methodology proposed herein will provide reasonable estimates of safe separation distances for the ranges of diameters, incident operating pressures and values of heat flux that have been previously identified.

## APPENDIX A - HEAT FLUX LEVELS

# Table A. 1 Summary of Heat Flux Levels 

| Heat Flux |  | Reference |  |
| :---: | :---: | :---: | :---: |
| Btu/hr $\mathrm{ft}^{2}$ | ${\mathrm{kW} / \mathrm{m}^{2}}$ | Number | Consequence |
| $222 * *$ | 0.7 | 37 | Exposed skin reddens and burns on prolonged exposure. |
| 250-380** | 0.79-1.2 | 32,35 | Heat flux of solar radiation. |
| 300 | 0.95** | *** | 115 seconds to severe pain; 663 seconds to second degree bum. |
| $317 * *$ | 1 | 34,39 | Level which is just tolerable by a clothed person. Also, this is the heat flux of solar radiation on a clear, hot summer day. |
| 380** | 1.2 | 56 | Solar heat flux at noontime during the summer. |
| $500 *$ | 1.57* | 39 | Flux limit from a flare if personnel are required to work in the area. |
| $500 *$ | 1.58* | 35 | Value of heat flux for a flare where personnel are continuously exposed. |
| 500 | 1.58** | 36 | Level at which personnel should not be exposed to for more than 2 hours. |
| 500 | 1.6 | 39 | Level to produce minor discomfort. |
| 507** | 1.6 | 32,34,56 | Minimum heat flux to be painful. Also, no discomfort felt for long exposures to the thermal radiation. |
| 550-555** | 1.74-1.75 | 35,37 | Pain threshold occurs after 60) second exposure. |
| 600 | 1.89** | *** | 45 seconds to severe pain; 187 seconds to second degree burn. |
| 634** | 2.0 | 37 | Causes damage to PVC insulated cables. |

Table A. 1
Summary of Heat Flux Levels (Continued)

Btu/hr $\mathrm{ft}^{2}$ | Heat Flux |
| :--- |
| $666^{* *}$ |

# Table A. 1 <br> Summary of Heat Flux Levels <br> (Continued) 

| $\underline{\text { Btu/hr } \mathrm{ft}^{2}}$ | $\begin{aligned} & \text { lux } \\ & \mathrm{kW} / \mathrm{m}^{2} \end{aligned}$ | Reference <br> Number | Consequence |
| :---: | :---: | :---: | :---: |
| 1500 | 4.73 | $35,40,57$ | Pain threshold for a $15-16$ second exposure. After this time period, blistering and permanent skin damage can occur. Also, this is the heat flux in areas where emergency actions occuring for several minutes may be performed by personnel without shielding but with appropriate clothing. This is also the level which can be tolerated for an indefinite period of time by an active worker wearing a hard hat, long sleeve shirt and gloves. |
| 1,500 | 4.73** | 36 | Level at which personnel should not be exposed to for more than 2 minutes. |
| 1,585** | 5.0 | 37 | Pain threshold reached after a 15 second exposure. |
| 1,600 | $5.05^{* *}$ | *** | 13 seconds to severe pain and 40 seconds to second degree burn |
| 1,650 | 5.2** | 57 | Heat flux which can be tolerated up to 5 minutes by a person wearing a hard hat, long sleeve shirt and gloves. |
| 1,900 | $5.99^{* *}$ | *** | 11 seconds to severe pain and 30 seconds to second degree burn. |
| 1,902** | 6 | 58 | Bare skin blisters if this heat flux is experienced for 20 seconds |
| 2,000 | $6.31^{* *}$ | 36 | Level at which personnel should not be exposed to for more than 20 seconds. |

Table A. 1<br>Summary of Heat Flux Levels<br>(Continued)

| Heat Flux |  | Reference |  |
| :---: | :---: | :---: | :---: |
| $B+u / h r \mathrm{ft}^{2}$ | $\mathrm{kW} / \mathrm{m}^{2}$ | Number | Consequence |
| $\begin{aligned} & 2,000- \\ & 2,029^{*} \end{aligned}$ | 6.3-6.4 | 35,37,38 | Pain threshold reached in 8 seconds and blistering occurs in 20 seconds. This is the maximum tolerable heat flux for shor term human exposure. This is also the heat intensity in areas where emergency actions lasting up to 1 minute may be required by personnel without shielding but with appropriate clothing. |
| 2,200 | 6.94 | 35 | 9 second exposure results in the pain threshold. |
| 2,500 | 7.89** | *** | 7 seconds to severe pain and 20 seconds to second degree burn. |
| 2,536** | 8 | 39 | Death occurs within minutes unless appropriate shelter is found. |
| 3,000 | 9.46 | 35,57 | Pain threshold reached in 6 seconds. This is also the value of heat flux (if solar radiation is included) for a flare at any location to which people have access (such as at grade underneath the flare or a service platform of a nearby tower). In this instance exposure should be limited to a few seconds since the heat flux level is adequate for escape only. Furthermore, limited over-exposure of a person wearing hard hat, long sleeve shirt and gloves results in a skin reaction similar to a mild sumburn. |
| $3.011^{* *}$ | 9.5 | $32,34,37$ | Pain threshold occurs in $6-8$ seconds, and second degree burns occur after 20 sec . |
| 3,200 | 10.09 | *** | 5 seconds to severe pain and 14 seconds to second degree burn. |

# Table A. 1 <br> Summary of Heat Flux Levels <br> (Continued) 

| Heat Flux |  | Reference |  |
| :---: | :---: | :---: | :---: |
| Btu/hr $\mathrm{ft}^{2}$ | $\mathrm{kW} / \mathrm{m}^{2}$ | Number | Consequence |
| 3,700 | 11.67 | 35 | Pain threshold reached at 4 seconds. |
| 3,709** | 11.7 | 56 | Mechanical integrity of thin, partially insulated steel can be lost. |
| 3,800 | 11.99** | *** | 4 seconds to severe pain and 11 seconds to second degree burn. |
| 3,962** | 12.5 | 32,34,37 | Piloted ignition of wood exposed to this heat flux for a prolonged period. Also, plastic tubing melts. |
| 3,994** | 12.6 | $\begin{aligned} & 32,34,39, \\ & 56 \end{aligned}$ | Exposed, dry, unpainted wood can beignited at this heat flux. Piloted ignition of cotton can also occur. This is the recommended flux level for residential areas at a frequency of $10 \times 10^{-6} / \mathrm{yr}$. Flammable materials in buildings may be damaged after 1,000 seconds. |
| 4,000 | 12.5 | 39 | Suitable heat flux level for control rooms or workshops. |
| 4.755** | 15.0 | 37 | This is the heat flux limit of Class 2 building materials. |
| 5,000* | 15.77 | 35 | Heat intensity on plant structures and in locations where operators are not likely to be working and where shelter from radiant heat is available. |
| 5,072** | 16.0 | 37,39 | Severe burns occur after 5 seconds, and wood ignites spontaneously. |
| 6,300 | 19.87 | 35 | Pain threshold reached in 2 seconds. |

## Table A. 1 <br> Summary of Heat Flux Levels <br> (Continued)

| Heat Flux |  | Reference |  |
| :---: | :---: | :---: | :---: |
| $\underline{\text { Bu/hr } \mathrm{ft}^{2}}$ | ${\mathrm{kW} / \mathrm{m}^{2}}$ | Number | Consequence |
| 6,340** | 20 | 58 | A four second exposure causes blistering of bare skin. This is also the heat flux level at which $1 \%$ lethality occurs for a twenty second exposure. |
| 7,291** | 23.0 | 56 | Mechanical integrity of thin, uninsulated steel can be lost. This is also the recommended flux level for potentially hazardous adjacent installations at a frequency of $50 \times 10^{-6}$. |
| 7,500 | 23.7 | 35 | White rats experience burns on their bare skin in approximately 6 seconds. |
| 7,925** | 25.0 | 37 | Prolonged exposure at this heat flux ignites wood. |
| 9,510** | 30.0 | 37 | Heat flux limit for Class I building materials. |
| 9,985** | 31.5 | 40 | After a few minutes exposure to this heat flux, wooden buildings, paper, window drapes and trees ignite spontaneously. |
| 10,778** | 34 | 58 | Piloted ignition of cellulosic material for a twenty second exposure. |
| 11,095** | 35 | 58 | $50 \%$ lethality occurs for a twenty second exposure to this heat flux. |
| 11,887** | 37.5 | 32,34 | Process equipment becomes damaged. This is also the minimum energy required to ignite wood at indefinitely long exposures. |

Table A. 1
Summary of Heat Flux Levels
(Continued)

| Heat Flux |  | Reference |  |
| :---: | :---: | :---: | :---: |
| Btu/hr $\mathrm{ft}^{2}$ | $\mathrm{kW} / \mathrm{m}^{2}$ | Number | Consequence |
| 11,983** | 37.8 | 34 | Flammable materials in process installations can be damaged after an exposure of 1,000 seconds. |
| 12,046** | 38 | 56 | Heat flux (for a 20 second exposure) corresponding to $50 \%$ mortality. |
| 20,922** | 66 | 58 | Piloted ignition of cellulosic material in four seconds. Also, $1 \%$ lethality for a four second exposure. |
| 37,089 | 117 | 58 | Heat flux resulting in $50 \%$ lethality for a four second exposure. |
| Various Flux Levels and Exposure Times |  | 39 | $(\mathrm{I}-25,400) \mathrm{t}^{4 / 5}=6,730$ <br> for spontaneous ignition of wood; 1 in units of $\mathrm{W} / \mathrm{m}^{2}$. |
|  |  | 39 | $(1-13,400) \mathrm{t}^{2 / 3}=8,050$ <br> for piloted ignition of wood; $I$ in units of $\mathrm{W} / \mathrm{m}^{2}$. |
|  |  | 39 | First degree burns assumed if $(t) 1^{1.15}$, 550,000 ; t in seconds and I expressed as $\mathrm{W} / \mathrm{m}^{2}$. |
|  |  | 34 | Graph of tolerance times to burn-injury levels for various incident heat fluxes. |

## Table A. 1 <br> Summary of Heat Flux Levels <br> (Continued)

| ${ }_{\text {Heat Flux }}{ }^{2}$ |  | Reference |  |
| :---: | :---: | :---: | :---: |
|  |  | Number | Consequence |
|  |  | 34 | Average heat flux which will cause severe blistering based on duration: $q_{2}=50 / t_{c}$ $q_{2}$ in $\mathrm{kW} /$ square meter and $\mathrm{t}_{\mathrm{c}}$ in seconds. |

NOTES

* Includes solar radiation
** Calculated using the relationship: $1 \mathrm{Btu} / \mathrm{hr} \mathrm{ft}^{2}=3.1546 \mathrm{~W} / \mathrm{m}^{2}$
*** Provided in material obtained from the TSB of Canada (personal communication. February 28, 1997).

APPENDIX B - BURN RADIUS CHARTS


Figure B.1-Burn Radius for 14" Diameter Pipeline $\mathrm{K}=3,962 \mathrm{Btu} / \mathrm{hr} \mathrm{sq} . \mathrm{ft}$.


Figure B. 2 - Burn Radius for 16 " Diameter Pipeline $\mathrm{K}=3,962 \mathrm{Btu} / \mathrm{hr} \mathrm{sq} . \mathrm{ft}$.


Figure B. 3 - Burn Radius for 18 " Diameter Pipeline $\mathrm{K}=3,962 \mathrm{Btu} / \mathrm{hr}$ sq. ft.


Figure B. 4 - Burn Radius for 20 " Diameter Pipeline $\mathrm{K}=3,962 \mathrm{Btu} / \mathrm{hr} \mathrm{sq} . \mathrm{ft}$.


Figure B. 5 - Burn Radius for 24" Diameter Pipeline $\mathrm{K}=3,962 \mathrm{Btu} / \mathrm{hr} \mathrm{sq} . \mathrm{ft}$.


Figure B.6-Burn Radius for 30" Diameter Pipeline $\mathrm{K}=3,962 \mathrm{Btu} / \mathrm{hr} \mathrm{sq} . \mathrm{ft}$.


Figure B. 7 - Burn Radius for 36" Diameter Pipeline $\mathrm{K}=3,962 \mathrm{Btu} / \mathrm{hr}$ sq. ft.


Figure B. 8 - Burn Radius for 14" Diameter Pipeline $\mathrm{K}=6,340 \mathrm{Btu} / \mathrm{hr}$ sq. ft.


Figure B. 9 - Burn Radius for 16" Diameter Pipeline $\mathrm{K}=6,340 \mathrm{Btu} / \mathrm{hr}$ sq. ft.


Figure B. 10 - Burn Radius for $18^{\prime \prime}$ Diameter Pipeline $\mathrm{K}=6,340 \mathrm{Btu} / \mathrm{hr} \mathrm{sq} . \mathrm{ft}$.


Figure B. 11 - Burn Radius for 20" Diameter Pipeline $\mathrm{K}=6,340 \mathrm{Btu} / \mathrm{hr}$ sq. ft .


Figure B. 12 - Burn Radius for 24" Diameter Pipeline $\mathrm{K}=6,340 \mathrm{Btu} / \mathrm{hr}$ sq. ft .


Figure B. 13 - Burn Radius for 30" Diameter Pipeline $\mathrm{K}=6,340 \mathrm{Btu} / \mathrm{hr} \mathrm{sq} . \mathrm{ft}$.


Figure B. 14 - Burn Radius for 36" Diameter Pipeline $\mathrm{K}=6,340 \mathrm{Btu} / \mathrm{hr}$ sq. ft .


Figure B. 15 - Burn Radius for 14" Diameter Pipeline $\mathrm{K}=9,510 \mathrm{Btu} / \mathrm{hr} \mathrm{sq} . \mathrm{ft}$.


Figure B. 16 - Burn Radius for 16" Diameter Pipeline $\mathrm{K}=9,510 \mathrm{Btu} / \mathrm{hr}$ sq. ft.


Figure B. 17 - Burn Radius for 18" Diameter Pipeline $\mathrm{K}=9,510 \mathrm{Btu} / \mathrm{hr}$ sq. ft .


Figure B. 18 - Burn Radius for 20" Diameter Pipeline $\mathrm{K}=9,510 \mathrm{Btu} / \mathrm{hr}$ sq. ft.


Figure B. 19 - Burn Radius for 24" Diameter Pipeline $\mathrm{K}=9,510 \mathrm{Btu} / \mathrm{hr} \mathrm{sq} . \mathrm{ft}$.


Figure B. 20 - Burn Radius for 30" Diameter Pipeline $\mathrm{K}=9,510 \mathrm{Btu} / \mathrm{hr}$ sq. ft.


Figure B. 21 - Burn Radius for 36" Diameter Pipeline $\mathrm{K}=9,510 \mathrm{Btu} / \mathrm{hr} \mathrm{sq} . \mathrm{ft}$.


Figure B. 22 - Burn Radius for 14" Diameter Pipeline $\mathrm{K}=9,985 \mathrm{Btu} / \mathrm{hr} \mathrm{sq} . \mathrm{ft}$.


Figure B. 23 - Burn Radius for 16" Diameter Pipeline $K=9,985 \mathrm{Btu} / \mathrm{hr} \mathrm{sq} . \mathrm{ft}$.


Figure B. 24 - Burn Radius for 18" Diameter Pipeline $\mathrm{K}=9,985 \mathrm{Btu} / \mathrm{hr} \mathrm{sq} . \mathrm{ft}$.


Figure B. 25 - Burn Radius for $20^{\prime \prime}$ Diameter Pipeline $\mathrm{K}=9,985 \mathrm{Btu} / \mathrm{hr}$ sq. ft.


Figure B. 26 - Burn Radius for 24" Diameter Pipeline $\mathrm{K}=9,985 \mathrm{Btu} / \mathrm{hr} \mathrm{sq} . \mathrm{ft}$.


Figure B. 27 - Burn Radius for 30" Diameter Pipeline $\mathrm{K}=9,985 \mathrm{Btu} / \mathrm{hr} \mathrm{sq} . \mathrm{ft}$.


Figure B. 28 - Burn Radius for $36^{\prime \prime}$ Diameter Pipeline $K=9,985 \mathrm{Btu} / \mathrm{hr}$ sq. ft.

## APPENDIX C - FIGURES AND TABLES FOR CHAPTER 5


Figure C.l - Detailed View of Edison Accident

## DURHAM WOODS



Figure C. 2 - Identification of Apartment Buildings


Figure C. 3 - Illustration of Apartments Used for Development of an Appropriate Scale
30000

Figure C. 4 - Comparison of Flux Levels Experienced Versus Levels for Blistering


| Diameter |
| :--- |
| Ratio |
| Selected |
| $16^{\prime \prime} / 14^{\prime \prime}$ |
| $18^{\prime \prime} / 14^{\prime \prime}$ |
| $24^{\prime \prime} / 14^{\prime \prime}$ |
| $30^{\prime \prime} / 14^{\prime \prime}$ |
| $36^{\prime \prime} / 14^{\prime \prime}$ |
| $18^{\prime \prime} / 16^{\prime \prime}$ |
| $24^{\prime \prime} / 16^{\prime \prime}$ |
| $30^{\prime \prime} / 16^{\prime \prime}$ |
| $36^{\prime \prime} / 16^{\prime \prime}$ |
| $24^{\prime \prime} / 18^{\prime \prime}$ |
| $30^{\prime \prime} / 18^{\prime \prime}$ |
| $36^{\prime \prime} / 18^{\prime \prime}$ |
| $30^{\prime \prime} / 24^{\prime \prime}$ |
| $36^{\prime \prime} / 24^{\prime \prime}$ |
| $36^{\prime \prime} / 30^{\prime \prime}$ |




APPENDIX D - VARIATION OF BURN RADIUS INFORMATION

The following pages provide tables, referenced in Chapter 6, which serve to demonstrate the influence of the following factors on the burn radius:

- Atmospheric transmissivity;
- Heat flux and atmospheric transmissivity (combined); and
- Wind direction and speed.

It should be remembered that the various burn radii are calculated using Equation 4-12.
Table D. 2
Percent Difference Between Burn Radii Obtained Using Assumed Vaue of $\tau$ and Burn Radii Determinued
Using Maximum and Minimum Values of $\tau$

Table D. 3
Variation of Burn Radii
For Combined Variations in Heat Flux and Atmospheric Transmissivity
Pipeline Conditions - Diameter/Incident Operating Pressure

| $36^{\prime \prime} / 1,200$ psia |  |  |
| :--- | :--- | :--- |
| Value of $\tau$ |  |  |
| $\frac{0,875}{1,559}$ | $\frac{0,746}{1,437}$ | $\frac{0,470}{1,132}$ |
| 1,345 | 1,239 | 975 |
| 1,199 | 1,104 | 866 |
| 1,225 | 1,128 | 885 |
| 1,055 | 970 | 758 |
| 938 | 862 | 671 |
|  |  |  |
| 992 | 912 | 711 |
| 852 | 782 | 606 |
| 756 | 692 | 533 |
|  |  |  |
| 967 | 889 | 692 |
| 830 | 762 | 589 |
| 736 | 674 | 518 |

36"/575 psia

 \begin{tabular}{lll}
\multicolumn{4}{c}{ Value of $\tau$} <br>
$\underline{0.875}$ \& $\underline{0.746}$ \& $\frac{0.470}{1,067}$ <br>
982 \& 768 <br>
918 \& 843 \& 656 <br>
815 \& 747 \& 578 <br>
833 \& 764 \& 592 <br>
713 \& 653 \& 500 <br>
630 \& 575 \& 436 <br>
\& \& <br>
668 \& 611 \& 466 <br>
568 \& 517 \& 388 <br>
498 \& 452 \& 333 <br>
\& \& <br>
650 \& 594 \& 452 <br>
552 \& 503 \& 376 <br>
484 \& 439 \& 321

 

\multicolumn{4}{c}{ Value of $\tau$} <br>
$\underline{0.875}$ \& $\underline{0.746}$ \& $\frac{0.470}{1,067}$ <br>
982 \& 768 <br>
918 \& 843 \& 656 <br>
815 \& 747 \& 578 <br>
833 \& 764 \& 592 <br>
713 \& 653 \& 500 <br>
630 \& 575 \& 436 <br>
\& \& <br>
668 \& 611 \& 466 <br>
568 \& 517 \& 388 <br>
498 \& 452 \& 333 <br>
\& \& <br>
650 \& 594 \& 452 <br>
552 \& 503 \& 376 <br>
484 \& 439 \& 321
\end{tabular} $\underset{+}{+}$ are varied.

Heat Flux
$\begin{aligned} & \text { Btu/hr } \mathrm{ft}^{2}\end{aligned}$
2,972
3,962
4,953
4,755
6,340
7,925
7,133
9,510
11,888
7,489
9,985
12,481
Note: Entries designated by "--." indicate conditions at which the burn radius is considered to be the reference value.

$$
\begin{aligned}
& \text { Equation 4-12. } \\
& 476 \\
& 525 \\
& 592 \\
& 689
\end{aligned}
$$


$\begin{aligned} \text { Note: } & \text { U.W. }=\text { Upwind } \\ & \text { D.W. }=\text { Downwind } \\ & \text { C.W. }=\text { Crosswind }\end{aligned}$
 Table D. 7
Variation of Burn Radius with Wind Direction and Speed
$36^{\prime \prime}$ Diameter Pipeline, $\mathrm{P}=1,000 \mathrm{psig}$
Distance to Specified Heat Flux - Feet

$9841,112 \quad 1,030$ ,

Note: $\begin{aligned} \text { U.W. }=\text { Upwind } \\ \text { D.W. }=\text { Downwind } \\ \text { C.W. }=\text { Crosswind }\end{aligned}$

## APPENDIX E - CANADIAN HEAT FLUX INFORMATION

## Exhibit No

$\qquad$

Undertaking given by Mr. Rothwell to Mr. Abes at Transcript<br>Volume 4, Page 468

Q. To provide a revised Response to NEB IR. No. 23 re Thermal

Radlation Curves.
A. The attached families of curves show the variation of thermal radaton flux with distance for NPS 24, 36 and 42 pipe at an operating pressure of 1000 psig and with wind speeds of 10 and $20 \mathrm{~km} / \mathrm{h}$. Other conditions a.e the same as those for Figures 23.2 to 23.9 of the original response. For each case, plots for upwind, downwind and crosswind directions are given.

Since Board Staff's request concerned the predictions immediately after ignition, we have elected not to incur the delay which would have been associated with the production of colour copies of the plots. However, 10 follow the evolution with tima of the thermal radiation flux on the monochrome copies, it is sufficient to recall that the curves are sequential with time from top to bottom, i.e the curve for time 0 is at the :op and the curve for :ime 900 is at the botom.

As was discussed qualitatively at Transcription Page 465, the distances :o specified lavels of thermal radiation increase somewhat in the downwind direction and decrease in the upwind direction. It is normal in cetaled risk assessments to take into account the statistical variation of wind cirection.












## REFERENCES

1. Transportation Research Board. Special Report 219 Pipelines and Public Safety. Washington, DC: Transportation Research Board.
2. National Transportation Safety Board. Texas Eastern Gas Pipeline Company Ruptures and Fires at Beaumont, Kentucky on April 27. 1985 and Lancaster. Kentucky on February 21, 1986. Washington, DC: National Transportation Safety Board (1987).
3. National Transportation Safety Board. Texas Eastern Transmission Corporation Natural Gas Pipeline Explosion and Fire Edison, New Jersey March 23, 1994. Washington, DC: National Transportation Safety Board (1995).
4. U.S. Department of Transportation, Research and Special Programs Administration. Pipeline Safety Regulations. Washington, DC: U.S. Department of Transportation (1993).
5. Griffis, F.H., R. Dresnack, S. Evans, E. Golub, J. Greenfeld, L.J. Pignataro, and L. Vega. Pipeline Industry: Comparison of U.S. with Foreign Pipeline Land Use and Siting Standards Maintenance, Rehabilitation and Retrofiting Policies and Practices. Final report prepared by NIIT under Contract No. DTRS 56-94-C-0006 for the U.S. Department of Transportation (April 1996).
6. Health \& Safety Executive. Risk Criteria for Land-Use Planning in the Vicinity of Major Industrial Hazards. London, England: Her Majesty's Stationery Office (1989).
7. Groves, T.K., P.R. Bishnoi, and J.M.E. Wallbridge. "Decompression Wave Velocities in Natural Gases in Pipe Lines." The Canadian Journal of Chemical Engineering 56 (1978): 664-668.
8. McAllister, E.W. Pipe Line Rules of Thumb Handbook. Houston, TX: Gulf Publishing Company (1993).
9. Picard, D.J., and P.R. Bishnoi. "The Importance of Real-Fluid Behavior in Predicting Release Rates Resulting from High-Pressure Sour-Gas Pipeline Ruptures." The Canadian Journal of Chemical Engineering 67 (1989): 3-9.
10. Vennard, J.K., and R.L. Street. Elementary Fluid Mechanics. New York, NY': John Wiley \& Sons (1982).

## REFERENCES

## (Continued)

11. Wilson, D.J. The Release and Dispersion of Gas from Pipeline Ruptures. Edmonton, Alberta, Canada: Alberta Environment (1979).
12. Wilson, D.J. Expansion and Plume Rise of Gas Jets from High Pressure Pipeline Ruptures. Edmonton, Alberta, Canada: Alberta Environment (1981).
13. Wilson, D.J. Predicting Risk of Exposure to Peak Concentrations in Fluctuating Plumes. Edmonton, Alberta, Canada: Alberta Environment (1982).
14. Fannelop, T.K., and I.L. Ryhming. "Massive Release of Gas from Long Pipelines." J. Energy 6 (1982): 132-140.
15. Flatt, R. "A Singly-Iterative Second-Order Method of Characteristics for Unsteady Compressible One-Dimensional Flows." Communications in Applied Numerical Methods 1 (1985): 269-274.
16. ----. "Unsteady Compressible Flow in Long Pipelines Following a Rupture," International Journal for Numerical Methods in Fluids 6 (1986): 83-100.
17. Lang, E., and T.K. Fannelop. "Efficient Computation of the Pipeline Break Problem." 3rd Symposium on Fluid Transients in Fluid Siructure Interaction. ASME 56 (1987): 115-123.
18. Ryhming, I. "On the Expansion Wave Problem in a Long Pipe with Wall Friction." Journal of Applied Mathematics and Physics 38 (1987): 378-390.
19. Olorummaiye, J.A., and N.E. Imide. "Computation of Natural Gas Pipeline Rupture Problems Using the Method of Characteristics." Journal of Hazardous Materials. 34 (1993): 81-98.
20. National Fire Protection Association and Society of Fire Protection Engineers. SFPE Handbook of Fire Protection Engineering. Quincy, MA: National Fire Protection Association (1988).
21. Hawthome, W.R., D.S. Weddell, and H.C. Hottell. "Mixing and Combustion in Turbulent Gas Jets." Third Symposium on Combustion, Flame and Explosion Phenomena. Combustion Institute (1949): 266-288.

## REFERENCES

## (Continued)

22. Gollahalli, S.R., T.A. Brzustowski, and H.F. Sullivan. "Characteristics of a Turbulent Propane Diffusion Flame in a Cross-Wind." Transactions of the CSME 3 (1975): 205-214.
23. Kalghatgi, G.T. "The Visible Shape and Size of a Turbulent Hydrocarbon Jet Diffusion Flame in a Cross-wind." Combustion and Flame 52 (1983): 91-106.
24. Oenbring, P.R., and T.R. Sifferman. "Flare Design...Are Current Methods Too Conservative?" Hydrocarbon Processing (May 1980): 124-129.
25. Fumarola, G., D.M. De Faveri, R. Pastorino, and G. Ferraiolo. "Determining Safety Zones for Exposure to Flare Radiation." I. Chem. E. Symposium Series No. 82 (1983): G23-G30.
26. U.S. Environmental Protection Agency and the National Oceanic and Atmospheric Administration. ALOHA User's Manual. Washington, DC: U.S. Environmental Protection Agency (1992).
27. Bakke, J.R., Dag Bjerketvedt, and Magne Bjorkhaug. "FLACS as a Tool for Safe Design Against Accidental Gas Explosions." I. Chem. E. Symposium Series No. 122 (1990): 141-152.
28. Raj, P.K., and J.A. Morris. "Computerized Spill Hazard Evaluation Models." Journal of Hazardous Materials 25 (1990): 77-92.
29. Drivas, P.J., J.S. Sabnis, and L.H. Teuscher. "Model Simulates Pipeline, Tank-Storage Failures." Oil \& Gas Journal (September 12, 1983): 162-169.
30. Technica International Ltd. WHAZAN User Guide. March 1988.
31. Hill, R.T., and J.R. Catmur. "Risks from Hazardous Pipelines in the United Kingdom." HSE Contract Research Report No. 82/1994 (1995).

32 Hockey, S.M., and P.J. Rew. "Review of Human Response to Thermal Radiation." HSE Contract Research Report No. 97/1996 (1996).
33. Golub. E., J. Greenfeld, R. Dresnack, F.H. Griffis, and L.J. Pignataro. Pipeline Accident Consequences for Natural Gas and Hazardous Liquid Transmission. Draft report prepared by NJIT under Contract No. DTRS 56-94-C-0006 for the U.S. Department of Transportation (November 1995).

## REFERENCES <br> (Continued)

34. Center for Chemical Process Safety of the American Institute of Chemical Engineers. Guidelines for Evaluating the Characteristics of Vapor Cloud Explosions, Flash Fires, and BLEVEs. New York, NY: American Institute of Chemical Engineers (1994).
35. American Petroleum Institute. Guide for Pressure-Relieving and Depressuring Systems. API Recommended Practice 521. Washington, DC: American Petroleum Institute (1990).
36. Brzustowski, T.A., and E.C. Sommer, Jr. "Predicting Radiant Heating from Flares." From Proceedings - Division of Refining. Washington, DC: American Petroleum Institute 53 (1973): 865-893.
37. Technica International Ltd. WHAZAN Theory Manual. March 1988.
38. Crawley, F.K. "The Effects of the Ignition of a Major Fuel Spillage." I. Chem. E. Symposium Series No. 71 (1982): 125-145.
39. Lees, F.P. Loss Prevention in the Process Industries, Boston, MA: Butterworth Inc. (1980).
40. Department of Housing and Urban Development. Safety Considerations in Siting Housing Projects. Washington, DC: Department of Housing and Urban Development (December 1975).
41. Transportation Safety Board of Canada. Transcanada Pipelines Limited Line 300-1, Natural Gas Pipeline Rupture Kilometre Post MLV 302-1 +2.849 KM Marionville, Ontario 1705 EDT, 06 June 1990. Hull, Quebec, Canada: Transportation Safety Board of Canada (1992).
42. National Transportation Safety Board. Northern Natural Gas Company Pipeline Puncture, Explosion and Fire Hudson, Iowa November 4, 1982. Washington, DC: National Transportation Safety Board (1983).
43. Transportation Safety Board of Canada. Transcanada Pipeliness Limited Natural Gas Pipeline Ruptures Main Line Valve $144-1+19.261 \mathrm{~km}$ Cardinat, Ontario 08 December 1991. Hull, Quebec, Canada: Transportation Safety Board of Canada (1994).

## REFERENCES <br> (Continued)

44. National Transportation Safety Board. Factual Pipeline Accident Report Basic Report (AR-01) FTW-79-F-P006 United Texas Transmission Co. Washington, DC: National Transportation Safety Board (1979).
45. Transportation Safety Board of Canada. Westcoast Energy Inc. 30-Inch Fort Nelson Main Line Natural Gas Pipeline Rupture Mile Post 102.6 Pink Mountain, British Columbia 1545 PDT, 06 October 1990. Hull, Quebec, Canada: Transportation Safety Board of Canada (1993).
46. Transportation Safety Board of Canada. Transcanada Pipelines Limited Natural Gas Pipeline Rupture Line 100-2, 914-Millimetre (36-Inch) Mainline Kilometre Post Main Line Valve 110-2 +22.098 Kilometres Latchford. Ontario 23 July 1994. Hull, Quebec, Canada: Transportation Safety Board of Canada (1995).
47. Transportation Safety Board of Canada. Natural Gas Pipeline Rupture Foothills Pipe Lines (Sask.) Lid. 1,067-Millimetre (42-Inch) Eastern Main Line Kilometre. Post $66+041$ Maple Creek, Saskatchewan 15 February 1994. Hull, Quebec, Canada: Transportation Safety Board of Canada (1995).
48. Felder, R.M., and R.W. Rousseau. Elementary Principles of Chemical Processes. New York, NY: John Wiley \& Sons (1978).
49. National Transportation Safety Board. Arkla Energy Resources Pipeline Rupture \& Fire Near Cale, Arkansas 02-24-86. Washington, DC: National Transportation Safety Board accident file.
50. National Transportation Safety Board. United Gas Pipe Line Company 20-Inch Pipeline Rupture and Fire Cartwright, Louisiana August 9, 1976. Washington, DC: National Transportation Safety Board (1977).
51. National Transportation Safety Board. Mobile Oil Corporation High-Pressure Natural Gas Pipeline, Near Houston, Texas September 9, 1969. Washington, DC: Nationa! Transportation Safety Board (1971).
52. National Transportation Safety Board. In the Matter of the Investigation of the Rupture and Fire Involving the Texas Eastern Gas Pipeline Company in Lancaster, Kentucky, on February 21, 1986. Washington, DC: from the National Transportation Safety Board accident file (1986).

## REFERENCES

## (Continued)

53. National Transportation Safety Board. Transcontinental Gas Pipe Line Corporation, 30-Inch Gas Transmission Pipeline Failure Near Bealeton, Virginia, June 9, 1974. Washington, DC: National Transportation Safety Board (1975).
54. National Transportation Safety Board. Pipeline Accident/Incident Summary Reports. Washington, DC: National Transportation Safety Board (1986).
55. Bair, F.E. The Weather Almanac. Detroit, MI: Gale Research Inc. (1992).
56. Bagster, D.F., and R.M. Pitblado. "Thermal Hazards in the Process Industry." Chemical Engineering Progress (July 1989): 69-75.
57. Schwartz, R., and M. Keller. "Environmental Factors Vs. Flare Application." CEP (September 1977): 41-44.
58. Roberts, A.F. "The Effect of Conditions Prior to Loss of Containment on Fireball Behavior." I. Chem. E. Symposium Series No. 71 (1982): 181-190.

[^0]:    * Converted from meters using the relationship 1 meter equals 3.2808 feet.

